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Multiple-reaction geobarometry for olivine-bearing igneous rocks

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Abstract

10

Efforts to map the vertical distribution of mafic and ultramafic igneous rocks in the Earth's crust and

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uppermost mantle have long been hampered by the lack of precise geobarometers for the appropriate

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mineral assemblages. The average pressure (avP) method (Powell and Holland 1994, *Am Min* 79, 120-

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133) is a multiple-reaction approach that uses a least-squares minimization to average the pressures

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derived from individual mineral equilibria, taking into account both their uncertainties and correlations.

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We applied avP to a carefully selected database of published phase equilibrium experiments in dry to

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H₂O-saturated, andesitic to basaltic and peridotitic systems at $P = 0.6\text{--}9.3$ kbar, $T = 940\text{--}1240^\circ\text{C}$, with

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$\log f\text{O}_2$ from NNO - 2.6 to NNO + 3.6 log units (where NNO is nickel-nickel oxide buffer). We made

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minor modifications to the thermodynamic models of clinopyroxene, spinel and olivine in order to

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improve the accuracy and precision of the results given by the avP method. Tests on the experimental

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database, using the modified thermodynamic models and spinel + clinopyroxene + olivine +

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plagioclase equilibria, showed that avP can reproduce the experimental P , within the calculated 1σ

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uncertainties (0.9–2.6 kbar; 1.6 kbar on average), for 67% of the database. No systematic deviations of

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the calculated pressure (P) with temperature (T) or mineral compositions are observed. Given the large

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compositional range of the experimental database, these results suggest that the method can be applied

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25 to any gabbroic, pyroxenitic or peridotitic rocks that contain the appropriate phase assemblage
26 clinopyroxene + olivine + plagioclase \pm spinel. For assemblages equilibrated at $P < 5$ kbar, the
27 calculated P shows a slight dependence on T , which therefore needs to be well constrained in order to
28 keep the overall P uncertainties as low as possible. T can be estimated using either available
29 independently-calibrated geothermometers or a simple calculation routine suggested in this work.
30 Application of avP to gabbroic xenoliths from Dominica, Lesser Antilles, and to gabbroic and
31 peridotitic xenoliths from Wikieup, Arizona, demonstrates the ability of the method to produce precise
32 P estimates for natural assemblages equilibrated at both mid and lower crustal conditions respectively.
33 Depending on the errors on mineral composition, appropriateness of the T estimate and attainment of
34 equilibrium of the assemblage, P uncertainty for natural rocks is ≤ 1.0 kbar. Such level of precision can
35 help discriminating rival petrogenetic processes in subduction zone, intra-plate and mid oceanic ridge
36 settings.

37 **Keywords**

38 Geobarometry, Phase equilibria, Mafic crust, Gabbros, Peridotites

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40 **Introduction**

41 Estimating the pressure (P) and temperature (T) of formation of mafic and ultramafic igneous rocks is a
42 fundamental step in studying the evolution of magmatic systems and the thermal and lithological
43 structure of the crust and crust-mantle transition zone. Developing accurate and precise
44 thermobarometric methods for such rocks is therefore critically important. Conventional methods use
45 phase compositions combined in either empirical or theoretical formulations to derive expressions for P
46 and T , with the fundamental assumption that the assemblage of interest attained equilibrium.
47 Commonly used geothermometers and geobarometers for gabbroic and peridotitic rocks are based on

48 single reactions between mineral end-members, often derived by simple regressions of experimental
49 data. Examples include two-pyroxene thermometers (e.g., Wells 1977; Putirka 2008) and barometers
50 (Putirka 2008), olivine-pyroxenes-plagioclase barometers (e.g., Fumagalli et al. 2017), hornblende-
51 plagioclase thermobarometers (e.g., Holland and Blundy 1994; Molina et al. 2015) and magnetite-
52 ilmenite thermometers (e.g., Ghiorso and Evans 2008). Empirical methods based on the composition of
53 a single mineral phase, such as clinopyroxene barometers (e.g., Nimis 1995; Putirka 2008) and
54 amphibole thermobarometers (e.g., Ridolfi and Renzulli 2012) are also available for gabbroic as well
55 as, basaltic and andesitic rocks. Mineral-melt thermobarometers potentially produce precise pressure
56 and temperature estimates (Falloon et al. 2007; Putirka 2008; Putirka 2016a; Neave and Putirka 2017),
57 but their application to gabbroic or peridotitic rocks is limited to assemblages containing melt in
58 equilibrium with the solid phases (interstitial melt or trapped melt inclusions).

59 For any given plutonic suite, a barometer or thermometer should meet the following conditions: i)
60 sufficient precision to discriminate between rival petrogenetic processes, ii) applicability over the P - T
61 range of interest, iii) calibration for the compositions of interest, iv) consistency across the entire suite.
62 Depending on the assemblage, criterion (i) can often be met for temperature, as model errors can be as
63 low as 40 °C (e.g., Putirka 2008), which is small relative to the range of formation of mafic/ultramafic
64 plutonic rocks (800–1300 °C). However, for existing barometers that are applicable to subsolidus
65 gabbroic assemblages (e.g., two-pyroxene or single-clinopyroxene barometers of Putirka 2008, single-
66 amphibole barometer of Ridolfi and Renzulli 2012), typical model errors are 3–4 kbar (inferred from
67 experimental data not used in the calibration; see Putirka 2008, 2016b and Erdmann et al. 2014 for a
68 review). This is unacceptably high considering the typical pressure range to be investigated (e.g., < 7
69 kbar for island arcs and intra-oceanic settings). For empirical or semi-empirical single-reaction
70 thermobarometers, condition (ii) is usually met, as calibrant phase equilibrium experiments generally
71 cover a wide range of P and T (e.g., 0–20 kbar, 800–1400°C). Condition (iii) is particularly critical
72 when adopting empirical or semi-empirical thermobarometry, because relationships that are too

73 simplified or that have no thermodynamic basis become unreliable when applied outside the range of
74 the calibration data. For the same reason, some empirical thermobarometers may be sensitive to the
75 buffering effect of phases that are not involved in the reaction, but are present in the assemblage,
76 hampering the application to natural samples with phase assemblages that differ from those in the
77 calibration dataset. This issue may be particularly important when single-phase empirical
78 thermobarometers calibrated using *suprasolidus* (i.e. melt-present) experiments (e.g., Ridolfi and
79 Renzulli 2012) are applied to plutonic rocks equilibrated at *subsolidus* (i.e. melt-absent) conditions.
80 Finally, condition (iv) may be difficult to meet when the rock suite includes samples with different
81 mineral assemblages, requiring the adoption of different, potentially inconsistent, single-reaction
82 thermobarometers calibrated on different experimental datasets.

83 A potential solution that fulfils all of the above conditions for pressure determination is multiple-
84 reaction barometry, employing an internally consistent dataset of thermodynamic data for mineral end-
85 members (e.g., Berman 1988; Holland and Powell 1998, 2011), along with a set of activity-
86 composition ($a-x$) relations calibrated for use with the same end-member dataset¹. The multiple-
87 reaction method allows all of the P and T information in a mineral assemblage to be used
88 simultaneously, provided the necessary thermodynamic models exist, via an independent set of
89 reactions among the end-members of the phases (Berman 1991; Gordon 1992; Powell and Holland
90 1988, 1994, 2008). End-member thermodynamic data and activity-composition ($a-x$) relations are now
91 available for the majority of rock-forming minerals of the crust and shallow mantle (e.g., Ghiorso and
92 Sack 1995; Holland and Powell 2011; Diener et al. 2007; Green et al. 2007, 2012, 2016; Gualda et al.
93 2012; White et al. 2014; Jennings and Holland 2015), thus potentially allowing the method to be
94 applied consistently to rock suites across a wide range of mineralogy and composition.

¹ We use the term ‘thermodynamic model’ to refer to both the thermodynamic properties of the end-members and the constituent $a-x$ relations

95 A single reaction among the end-members of phases is sufficient to calculate the value of P at given
96 T . However, in a robust application of multiple-reaction barometry, there are several reactions forming
97 an independent set, and the P values are statistically over-determined. Owing to the uncertainties in
98 thermodynamic models and in analysed mineral compositions, the reactions do not, in general, intersect
99 at the same $P(T)$. Therefore a statistical approach is needed to calculate the optimal P . Such an
100 approach must be weighted by the uncertainties affecting each reaction in the independent set, and this
101 weighting must also take into account the correlated influence of these uncertainties on reactions that
102 share end-members (Powell 1985; Powell and Holland 1994). The average P (avP) method (Powell and
103 Holland 1988, 1994, 2008), is one of the few available algorithms for multiple-reaction
104 thermobarometry that uses a rigorous statistical approach to handle both uncertainties *and* correlations.
105 The algorithm is accessible via the software THERMOCALC, using the Holland and Powell (1985, 1990,
106 1998, 2011) internally consistent dataset of end-members and, in principle, any suitably encoded a - x
107 relations with the modified regular solution formulations described in Powell and Holland (1993) and
108 Holland and Powell (2003).

109 The avP algorithm in THERMOCALC has historically been applied to metamorphic rocks, and the
110 development of a - x relations suitable for their use has concentrated on the metamorphic domain.
111 However, following the work of Jennings and Holland (2015) and Green et al. (2016), a set of a - x
112 relations is, for the first time, available for the phases expected in hydrous igneous rocks under crustal
113 and upper mantle conditions. Given the potential of avP to fulfil all the conditions for optimal
114 geobarometry, we evaluated the accuracy of the method and accompanying a - x relations for mafic and
115 ultramafic rocks using a dataset of phase equilibrium experiments in basaltic and peridotitic systems,
116 together with well-equilibrated natural samples, as a test case.

117 **Methodological overview and rationale**

118 Our aim is to develop a procedure for reliably applying multiple-reaction barometry to crustal-depth
119 mafic and ultramafic igneous rocks. A database of phase equilibrium experiments (Supplementary
120 Material 1) was used to evaluate the results obtained with THERMOCALC's avP algorithm (Powell and
121 Holland 1988, 1994, 2008) and a selection of published thermodynamic models for mineral phases.
122 The thermodynamic models were then refined, where necessary, to minimise systematic discrepancies
123 between the avP calculations and the experimental database. Prior estimates of T and analytical or
124 model uncertainty were explored, both for the experimental database, and for natural samples. Our
125 work focused on rocks containing the assemblage SCOIP (see Table 1 for abbreviations). This
126 assemblage is common in mafic and ultramafic plutonic rocks from a variety of tectonic settings, and
127 appears in a relatively large number of experimental runs. We further tested avP calculations using the
128 sub-assemblage COIP. This yields only three independent reactions, compared with six for SCOIP, and
129 is therefore potentially less robust. However, COIP can be applied to a larger number of experiments
130 and natural samples, including plagioclase peridotites.

131 Refinement of the thermodynamic models allows the avP approach to satisfy criteria (i) to (iii)
132 above. Regarding criterion (iv), the method is devised such that it is possible to apply mutually
133 consistent barometers across a suite of rocks, simply by adding further sets of well-calibrated $a-x$
134 relations to those for olivine, clinopyroxene, plagioclase and spinel, using the same internally-
135 consistent dataset of end-member thermodynamic properties (Holland and Powell 2011).

136 Supplementary Material 2 describes how to obtain software and input files for applying the avP
137 approach to SCOIP and COIP assemblages, and provides an example calculation.

138 **Experimental database**

139 We assembled a database of published phase equilibrium experiments in natural andesitic to basaltic
140 and peridotitic systems whose run products contain, at least, olivine, clinopyroxene and plagioclase.
141 Some experimental datasets were excluded, despite containing the right assemblage, either because (i)

142 the experiments show clear evidence of disequilibrium (e.g., large compositional inhomogeneity of the
143 run products, significant Fe loss), or (ii) run duration was very short (i.e. < 2 hours) or (iii) mineral
144 compositions were not reported. A reference list of experiments that satisfied these requirements is
145 reported in Supplementary Material 1 and their P - T conditions are shown in Figure 1. Clinopyroxenes
146 in most of the experiments at $P \geq 9$ kbar and $T \geq 1150^\circ\text{C}$ have low Ca contents (≥ 0.38 atoms per
147 formula unit, apfu; Bartels et al. 1991; Draper and Johnston 1992; Kinzler and Grove 1992; Fig. 2a).
148 The closure of the solvus separating low- and high-Ca pyroxene compositions is difficult to capture
149 accurately in a - x modelling and, moreover, such low-Ca clinopyroxenes are rare in natural igneous
150 rocks. We therefore excluded the experiments containing these clinopyroxenes from the dataset.
151 Additionally, clinopyroxenes from atmospheric pressure experiments (0.001 kbar) show a range of
152 tetrahedral aluminium ($\text{Al}^{\text{IV}} = 2 - \text{Si}$ [apfu]) contents that exceeds that at higher pressures (Fig. 3),
153 suggestive of a complex cation partitioning for 1-atm clinopyroxenes. Existing geobarometers
154 involving Al- and Na-in-clinopyroxene components systematically fail to predict pressure for 1-atm
155 experiments. This has been ascribed to Na-loss during the experiment (Putirka 2008), although,
156 conversely, Mollo et al. (2010) showed that fast-growing clinopyroxene crystals contain excessive Na
157 and Al, compensating for depletions in Si, Ca and Mg relative to the equilibrium composition.
158 Considering that some equilibria used in this work involve Al and Na partitioning, we therefore
159 removed all 1-atm experiments from the avP dataset. Additional filtering based on the quality of
160 microprobe analyses and possible presence of aberrant data (see Supplementary Material 1) yielded a
161 final dataset of 209 experiments with the COIP assemblage, of which 62 also contain the SCOIP
162 assemblage.

163 The conditions covered by the refined experimental database are $P = 0.5$ – 9.3 kbar, $T = 935$ – 1240°C
164 (Fig. 1) and $\log f\text{O}_2$ from NNO - 2.6 to NNO + 3.6 log units. Compositions of olivine, spinel,
165 clinopyroxene and plagioclase are shown in Fig. 2. Olivines are almost pure solid solutions of forsterite

166 and fayalite, with only minor or trace amounts of Ca and Mn components ($\text{Ca} < 0.025$ apfu; $\text{Mn} <$
167 0.024 apfu). Plagioclases are solid solutions of albite and anorthite, with only minor amounts of K ($<$
168 0.12 apfu) and Fe_{tot} (< 0.09 apfu) components. Spinel shows a wide range of composition including
169 magnetites, ulvospinel and spinel *sensu stricto*, with variable amounts of Cr (Fig. 2c) and mg# in the
170 range 0.11–0.89. Clinopyroxenes are augites with variable Al (0.08–0.39 apfu) and Ti (0.00–0.10 apfu)
171 and low Na (0.01–0.08 apfu). The 62 experiments containing the subset SCOIP cover roughly the same
172 range of P–T conditions and mineral compositions (e.g., Figs. 1 and 2a,b), with $\log f\text{O}_2$ from NNO -0.1
173 to NNO + 3.6.

174 **Average-Pressure calculations: methodology**

175 **Overview of the avP method**

176 An optimal value of P , \bar{P} , may be found by least-squares averaging of multiple independent phase
177 equilibrium constraints, each of which can be formulated as a reaction between end-members. To find
178 the optimal value is not simply a matter of averaging the values of the P_r , the pressures obtained from
179 each reaction r in the set, because some end-members appear in more than one reaction. If the pressures
180 of the reactions are adjusted to coincide at \bar{P} , the implication is that the calculated chemical potentials
181 of the end-members are being adjusted. Adjustments to chemical potentials must be applied
182 consistently across each of the end-member reactions, meaning that the reactions must be moved in a
183 suitably correlated way in P space in order to coincide at \bar{P} . The approach is most powerful if applied
184 via using a set of fully thermodynamic models for the phases, in which the chemical potentials of end-
185 members are constructed from parameters representing other physical properties, such as standard
186 enthalpies, within some uncertainty. In this case, the correlated adjustments to the P_r required to
187 achieve coincidence at \bar{P} , can be made by adjusting the end-member properties, based on a covariance
188 matrix that is known *a priori*. The method amounts to ‘least-squares adjustment of indirect

189 observations' (e.g. Mikhail 1976). Such a method, the avP algorithm, is described in detail by Powell
190 and Holland (1988, 1994, 2008), and exploits the Holland and Powell (1990, 1998, 2011) dataset of
191 end-member thermodynamic properties, combined with a set of a – x relations for each solid solution to
192 describe the thermodynamics of mixing (Gordon 1992 discussed a similar procedure). The TWEEQU
193 thermobarometric method of Berman (1991) is of the same type, but does not account for the
194 correlations in the adjustments applied to the P_r .

195 A statement of equilibrium for each multiphase reaction, r , can be written as

196
$$\sum_i \nu_i \mu_i = \sum_i \nu_i \mu_i^0 + R T \ln K_r = \Delta G_r^\circ + RT \ln K_r = 0, \quad (1)$$

197 where, for end-member i in a phase j of specified crystal structure, ν_i is the reaction coefficient, μ_i is
198 the total chemical potential and μ_i^0 is the chemical potential of the pure end-member at P , T . ΔG_r° is the
199 Gibbs free energy difference in kJ for the reaction among the pure end-members, R is the gas constant
200 (0.0083144 kJ/K·mole) and K_r is the equilibrium constant, which is given by

201
$$K_r = \prod_i (a_i)^{\nu_i}, \quad (2)$$

202 a_i being the activity of end-member i .

203 The avP algorithm, as implemented in THERMOCALC (Powell and Holland 1988), finds the activities
204 of end-members given the observed compositions of the phases, at a likely P_0 and T_0 (initial estimates
205 of pressure and temperature), to compute each equilibrium constant K_r . Taking the corresponding
206 expression for ΔG_r° from the Holland and Powell (1985 1990, 1998, 2011) dataset, it linearises
207 expression (1) at T_0 , P_0 , to form the equation from which P is obtained (Fig. 4):

208
$$a + b T_0 + c P + RT_0 \ln K_r = 0 \quad (3)$$

209 The a , b and c terms represent the end-member enthalpy, entropy and volume changes, ΔH_r , $-\Delta S_r$ and
210 ΔV_r , respectively, for the reaction.

211 Uncertainties in equation (3) are derived from uncertainties in the properties of individual end-
212 members involved in the reaction. Only the enthalpies and activities of the end-members are considered

213 to be uncertain, since the Holland and Powell dataset treats end-member entropies and volumetric
214 properties as exact. The uncertainties in the end-member enthalpies are correlated with each other, and
215 their standard deviations and correlation coefficients are extracted from the Holland and Powell dataset.
216 The uncertainties in the activities of end-members are assumed to be independent of each other, and are
217 propagated through from (i) the analytical uncertainties in the compositional variables and (ii) the
218 uncertainties in the a - x relations, treated simply as uncertainties in the values of $W_{(m,n)}^j$.

219 As described in Powell and Holland (1988), the uncertainties in the enthalpies and activities of end-
220 members can be converted into uncertainties on the a and $\ln K_r$ terms of equation (3), taking into
221 account the reaction coefficient of each end-member. A full covariance matrix can ultimately be
222 derived, describing the uncertainties and correlations among the estimates of P from each of the
223 independent reactions (P_r).

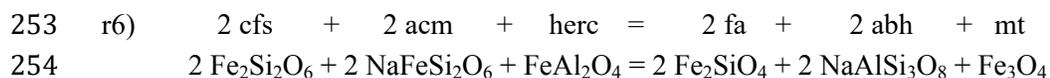
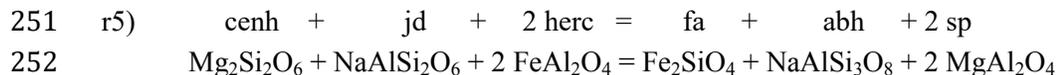
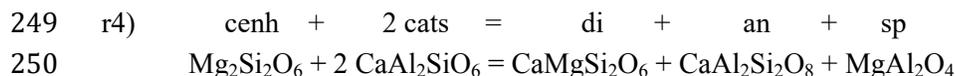
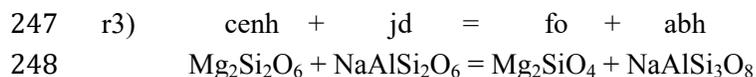
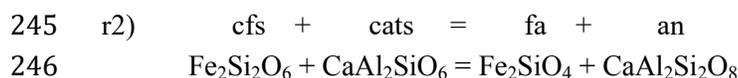
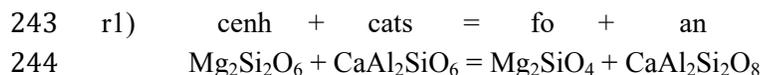
224 Using the reactions and the corresponding covariance matrix, the avP algorithm finds an optimal
225 least-squares average of the P_r values, \bar{P} , with a well-defined uncertainty $\sigma_{\bar{P}}(T_0)$. Figure C2 from
226 Powell and Holland (1988) illustrates graphically how two correlated pressures (P_r) are combined to
227 give \bar{P} .

228 A number of diagnostic values are displayed following the avP procedure (see Supplementary
229 Material 2 and Powell and Holland, 1994), to assist meaningful application of the method. The
230 diagnostic value σ_{fit} is particularly helpful, providing a measure of how appropriately the pressure
231 information in the independent set of reactions can be combined. If σ_{fit} lies within the cut-off value
232 provided by the χ^2 test, representing 95 % confidence, then it is reasonable to assume that the selected
233 phases are indeed in thermodynamic equilibrium, given realistic uncertainties in the mineral analyses
234 and thermodynamic models.

235 **Independent reactions used in this work**

236 For each of the two assemblages (SCOIP and COIP) considered, we specified the independent set of

237 reactions to be used by the avP algorithm. These were chosen to have generally low uncertainties in
238 their individual pressure estimations, P_r , implying that they are sensitive to pressure, and also to lie
239 reliably within 20 kbar of the experimental pressures during model refinement. The choice of
240 independent reactions affects the calculated \bar{P} , but the variation in \bar{P} is of the order of $\sigma_{\bar{P}}$. For the
241 assemblage SCOIP, the selected reactions are (see example in Supplementary Material 2 for ΔH_r , ΔS_r
242 and ΔV_r for each reaction):



255 where the number of independent reactions is equal to $s - c$, with s being the number of end-members
256 and c the number of components involved in mixing (see Table 2).

257 For COIP, an independent set contains three reactions. The natural choice is simply those of the set
258 of reactions r1—r6 that do not involve end-members of spinel solid solution, i.e., reactions r1, r2 and
259 r3.

260 **Thermodynamic models**

261 We employed the internally consistent thermodynamic dataset of Holland and Powell (2011; version
262 6.3, generated 29 March, 2016), and an initial selection of published $a-x$ relations that are appropriate

263 for use with this dataset: the models for plagioclase, clinopyroxene, olivine and spinel were taken
264 respectively from Holland and Powell (2003), Green et al. (2016), Hackler and Wood (1989), and
265 Bryndzia and Wood (1990). Specifically, these $a-x$ relations are regular solution models (one-
266 parameter Margules solutions), with mixing on sites, allowing for asymmetry introduced via a van Laar
267 parameter applied to each end-member.

268 During our development of a procedure for application of avP to igneous rocks, a major requirement
269 was to make modifications to the thermodynamic models (end-member properties and/or $a-x$
270 relations) in order to achieve accurate estimation of \bar{P} values for our experimental database. Many of
271 these modifications involved the adjustment of end-member Gibbs energy functions from the dataset,
272 $G^{ds}(P,T)$, by functions $\Delta G^{mod}(P,T)$, such that for end-member i , $G_i(P,T) = G_i^{ds}(P,T) + \Delta G_i^{mod}(P,T) =$
273 $G_i^{ds}(P,T) + a + b T + c P$ (the constants a , b and c may be considered to represent enthalpy, entropy
274 and volume terms respectively). Additionally we made minor structural changes to the $a-x$ relations,
275 and adjusted some of the interaction energies, $W(m,n)$.

276 By repeatedly assessing the results of avP run on the experimental database as described below, we
277 adopted the following thermodynamic models for clinopyroxene, olivine and spinel. We did not modify
278 the Holland and Powell (2003) models for plagioclase, an asymmetric solution between high-albite,
279 anorthite and sanidine. Model end-members are listed in Table 2, and a full description of the $a-x$
280 relations and end-member thermodynamic data can be found in the Supplementary Material 3.

281 **Clinopyroxene.** The starting point for the clinopyroxene $a-x$ relations was the ‘augite model’ of
282 Green et al. (2016) that entails mixing of [Mg, Fe²⁺, Al, Fe³⁺] on the M1, [Na, Ca, Mg, Fe²⁺] on the
283 M2, and [Si, Al] on the tetrahedral sites. The end-members high-temperature clinoenstatite (cenh) and
284 clinoferrosilite (cfs) do not appear in the Holland and Powell (2011) dataset, but are constructed by
285 adding $\Delta G_i^{trans}(P,T)$ terms to the Gibbs energy functions for ortho-enstatite (en) and ortho-ferrosilite
286 (fs) respectively (Green et al. 2016), where $\Delta G_i^{trans}(P,T)$ represents the Gibbs energy change of the

287 *Pbca*—*C2/c* transition. The model incorporates order-disorder of Fe²⁺-Mg across M1 and M2,
288 expressed via the proportions of an end-member *fmc*, Fe^{M2}Mg^{M1}Si₂O₆. Like the *cenh* and *cfs* end-
289 members, the *fmc* end-member does not appear in the dataset, but has a Gibbs energy function that can
290 be expressed as $\frac{1}{2} (G_{\text{cenh}}(P, T) + G_{\text{cfs}}(P, T)) + \Delta G_{\text{fmc}}^{\text{od}} = \frac{1}{2} (G_{\text{en}}^{\text{ds}}(P, T) + G_{\text{fs}}^{\text{ds}}(P, T) + \Delta G_{\text{en}}^{\text{trans}}(P, T) +$
291 $\Delta G_{\text{fs}}^{\text{trans}}(P, T)) + \Delta G_{\text{fmc}}^{\text{od}}$, where $\Delta G_{\text{fmc}}^{\text{od}}$ is the Gibbs energy change of ordering to form *fmc* from a
292 disordered 1:1 mixture of *cenh* and *cfs*. The model also includes Si-Al order-disorder on the tetrahedral
293 site, but whereas in Green et al. (2016) this order-disorder was expressed explicitly, by including both
294 an ordered and a disordered Ca-tschermaks end-member in the solid solution, we here made a minor
295 simplification by replacing the ordered and disordered Ca-tschermaks end-members with a single end-
296 member, *cats*, from the Holland and Powell (2011) dataset, for which the thermodynamic properties
297 reflect disordering with temperature. This makes negligible difference to the calculations.

298 Three problems that emerged during testing of the avP procedure were linked to the clinopyroxene
299 model:

- 300 1) Calculations on the COIP assemblage tended to produce \bar{P} estimates that exceeded experimental
301 pressures by ~5 kbar.
- 302 2) The avP diagnostics (see Supplementary Material 2) for the calculations on the SCOIP assemblage
303 suggested that the clinopyroxene end-members Ca-tschermaks and jadeite were particularly likely
304 to exhibit large, uncertainty-normalised residuals in activity, e^* , and that *cenh* strongly influenced
305 \bar{P} .
- 306 3) When \bar{P} was calculated over a range of T_θ , values of σ_{fit} and $\sigma_{\bar{P}}$ were typically smaller for small
307 values of temperature, with minima occurring at temperatures $< T_{\text{exp}}$. This showed that the models'
308 temperature-dependences were poorly calibrated. This is a significant observation since the
309 method of calculating \bar{P} at multiple T_θ may be used to estimate the equilibration temperature of a
310 sample, as discussed below.

311 In response to problem 3), we investigated the model reaction Mg_2SiO_4 (fo) + $\text{Fe}_2\text{Si}_2\text{O}_6$ (cfs) = Fe_2SiO_4
312 (fa) + $\text{Mg}_2\text{Si}_2\text{O}_6$ (cenh), which is not included explicitly among our reactions, but appears implicitly
313 since it is obtained by subtracting r2 from r1. Taking both our own database and that of Loucks (1996),
314 we recalibrated the clinopyroxene and olivine models, treating fo + cfs = fa + cenh as a single-reaction
315 thermometer. The key changes made to the clinopyroxene model through this approach were to
316 $\Delta G_{\text{cen}}^{\text{tran}}(P,T)$ and $\Delta G_{\text{cfs}}^{\text{tran}}(P,T)$, which were recalibrated while preserving their fit to the estimated
317 transition curves in P – T space (taken from Boyd et al. 1964, as interpreted by Gasparik 1990, for
318 enstatite-clinoenstatite; and from Lindsley 1981, for ferrosilite-clinoferrosilite). We also decreased the
319 value of $\Delta G_{\text{fmc}}^{\text{od}}(P,T)$ from -4.4 kJ mol^{-1} to -6.5 kJ mol^{-1} . Minor adjustments were made to the
320 relatively well-known interaction energies in the CaO–FeO–MgO–SiO₂ system, $W_{(\text{di},\text{cfs})}$, $W_{(\text{di},\text{fmc})}$,
321 $W_{(\text{cenh},\text{cfs})}$, and $W_{(\text{cfs},\text{fmc})}$, while respecting the experimental constraints placed on them (Turnock and
322 Lindsley 1981; Lindsley 1983; see Green et al. 2016). The revised values are considered to be better
323 constrained than the originals. The modified models continue to predict temperatures for the fo + cfs =
324 fa + cenh reaction with a large scatter, but now with a reasonable standard error of estimate of $\pm 71^\circ\text{C}$.

325 We made further adjustments to many of the less well-constrained interaction energies, with
326 reference to both the fo + cfs = fa + cenh reaction and the results of avP calculations on the
327 experimental database. These adjustments involved reducing the size of the interaction energies by
328 ~40%. In modifying the interaction energies, we treated them as falling into groups that should retain
329 similar values: (1) the $W([\text{cenh},\text{cfs},\text{fmc}],[\text{jd},\text{acm}])$ parameters (notation indicates interaction energies
330 between each of the end-members in the first square brackets with each of the end-members in the
331 second); (2) the $W([\text{cenh},\text{cfs},\text{fmc}],[\text{cats}])$ parameters; (3) the $W([\text{jd},\text{acm}],[\text{cats}])$ parameters.
332 $W([\text{di}],[\text{cats}])$ was adjusted independently. These modifications greatly reduced the excessive avP
333 estimates for COIP (Fig. 5b), and removed the high residuals associated with a_{cats} and a_{jd} when SCOIP
334 is used. Furthermore, values of σ_{fit} for the experiments typically decreased, with fewer experiments

335 failing the χ^2 test; this was a valuable indication that the adjustments tended to make the models, as a
336 set, predict a more mutually consistent set of independent reactions given the experimental phase
337 compositions.

338 Individual changes made to the poorly-constrained interaction energies should not necessarily be
339 regarded as improvements to the original $a-x$ relations, since avP in general demands less accurate $a-x$
340 relations than the forward-modelling calculations for which the Green et al. (2016) clinopyroxene
341 model was developed. However it is interesting that a much less non-ideal model was required to
342 produce good avP results, especially considering that the pressure information in reactions r1-r3 is
343 related primarily to formation of aluminous pyroxene from plagioclase (i.e. Ca-tschermak + Quartz =
344 Anorthite). The thermodynamics of the key pyroxene and plagioclase end-members in the Holland and
345 Powell (2011) dataset, including the crucial volumetric parameters, are expected to be well known.

346 **Olivine.** We initially followed Jennings and Holland (2015) in adopting a simple symmetric model for
347 a forsterite—fayalite solid solution, with Mg—Fe mixing on two identical M sites, as per Hackler and
348 Wood (1989), with $W(\text{fo,fa}) = 8 \text{ kJ mol}^{-1}$. During refitting of the reaction $\text{fo} + \text{cfs} = \text{fa} + \text{cenh}$, we found
349 it necessary to introduce order-disorder, via an ordered end-member $\text{Fe}^{\text{M1}}\text{Mg}^{\text{M2}}\text{SiO}_4$ (ol_{fm}; compare
350 with discussion of ordered end-member fmc above; see also Supplementary Materials 3), with a Gibbs
351 energy function given by $G_{\text{ol}_{\text{fm}}}(P,T) = \frac{1}{2} \{G_{\text{fo}}(P,T) + G_{\text{fa}}(P,T)\} + \Delta G_{\text{ol}_{\text{fm}}}^{\text{od}}(P,T)$, and $\Delta G_{\text{ol}_{\text{fm}}}^{\text{od}}(P,T) = (-8$
352 $+ 0.003 T + 0.02 P) \text{ kJ mol}^{-1}$. This expression yields -3.8 kJ mol^{-1} at 1100°C and 5 kbar, and remains
353 negative, i.e. favouring ordering, below 2300°C within the pressure range of interest. $\Delta G_{\text{ol}_{\text{fm}}}^{\text{od}}(P,T)$ is
354 thus of a physically plausible size over the P - T range of interest, comparable with $\Delta G_{\text{fmc}}^{\text{od}}(P,T)$. We
355 stress, however, that our calibration does not provide unique constraints on the a , b and c terms of
356 $\Delta G_{\text{ol}_{\text{fm}}}^{\text{od}}(P,T)$ and they should be considered essentially as fit parameters, compensating for other
357 difficulties in modelling Fe-Mg exchange between olivine and clinopyroxene in natural systems.

358 **Spinel.** Jennings and Holland (2015) used the spinel $a-x$ relations of Bryndzia and Wood (1990), in
359 the system MgO—FeO—Fe₂O₃—Al₂O₃—Cr₂O₃, in which all cations mix on three identical sites. Such
360 a model greatly simplifies the mixing behaviour of spinel, by neglecting cation and magnetic
361 disordering. Nevertheless, Bryndzia and Wood (1990) applied it successfully to oxy-barometry of
362 spinel peridotites, and it performed convincingly in the forward calculations of Jennings and Holland
363 (2015), also on peridotite. Therefore we initially adopted the Bryndzia and Wood (1990) $a-x$ relations,
364 extended to include TiO₂ by T.J.B. Holland (personal communication; see Supplementary Material 3).

365 For the majority of experiments that contained magnetite rather than Cr/Al-spinel, we found that avP
366 calculations with SCOIP had a tendency to underestimate experimental pressures, even when
367 calculations with COIP tended to overestimate them. Modifications to the clinopyroxene model, as
368 described above, reduced \bar{P} estimates for both assemblages. To overcome this discrepancy between the
369 assemblages, and following the observation that the end-members sp and herc showed large
370 uncertainty-normalised residuals in activity (as suggested by the avP diagnostics), we were obliged to
371 add large $\Delta G_i^{mod}(P,T)$ terms, of 5 kJ mol⁻¹, to the sp and herc end-members. These terms may be
372 crudely compensating for the highly simplified form of the spinel $a-x$ relations, and indeed any
373 unidentified weaknesses elsewhere in the set of thermodynamic models, that compensate for each other
374 in COIP. Additionally, we found that pressures for experiments containing Cr/Al-spinel were routinely
375 overestimated. We were able to prevent this by adjusting the value of $W_{sp-pier}$ from 39 to 23 kJ mol⁻¹,
376 and that of $W_{herc-pier}$ from 27 to 16 kJ mol⁻¹. The modifications also allowed most of the experiments to
377 pass the χ^2 test, whereas, using the original spinel model, most of SCOIP calculations failed, with σ_{fit}
378 $\gg 1.49$.

379 **Inputs for avP: compositions, uncertainties, P_0 and T_0**

380 Mineral formulae for the SCOIP/COIP phases in the experimental runs and natural samples were
381 recalculated from the compositions determined by electron microprobe. Iron was assumed to be all

382 ferrous in olivine and all ferric in plagioclase. Fe^{3+} in clinopyroxene and spinel was calculated from
383 stoichiometry following the method of Droop (1987). Mineral formulae were then converted into
384 compositional variables suitable for the $a-x$ relations (Supplementary Material 2). THERMOCALC
385 (Powell and Holland 1988) version 3.47i was used to run the avP algorithm on these data, generating a
386 value of \bar{P} based on each assemblage for each experiment. The results obtained from the avP algorithm
387 are sensitive to the relative size of uncertainties specified for the compositional variables, and also to
388 those applied to the interaction energies, $W_{(m,n)}$, in the $a-x$ relations.

389 Since rigorous error propagation is not applicable to our experimental database (see Supplementary
390 Material 4), uncertainties on most of the compositional variables for the experimental samples have
391 been simply set to ± 0.01 apfu, a default value typically adopted in avP (Powell and Holland 2008).
392 Given the difficulty of quantifying realistic uncertainties on the spinel $a-x$ model, which does not
393 account for cation ordering, uncertainties on $x^{\text{sp}} [\text{Fe}^{2+}/(\text{Fe}^{2+}+\text{Mg})]$ and $f^{\text{sp}} [\text{Fe}^{3+}]$ have been set to 0.04
394 apfu, with the aim of accounting for the possible oversimplification of spinel $a-x$ relations. This
395 simplified set of compositional uncertainties was then compared with the true uncertainties, derived
396 from partial error propagation, during trial calculations on the refined experimental database using
397 SCOIP equilibria. No significant difference was found in the results. In this test, we also noted that the
398 maximum variations of \bar{P} between different uncertainties assumptions are on the order of $\sigma_{\bar{P}}$, as
399 suggested by Powell and Holland (2008).

400 For the natural samples, the mineral compositions used for thermobarometry are taken from
401 individual spot analyses of adjacent crystal rims, since adjacent rims are assumed to most closely
402 represent the equilibrium assemblage. The compositional errors are initially derived from counting
403 statistics of the electron microprobe (Supplementary Table 5). Given that such errors are small and
404 weakly correlated, partial error propagation is assumed to yield sufficiently realistic uncertainties for
405 most of the $a-x$ model compositional variables. However, after applying the partial error, uncertainties

406 on compositional variables $x(\text{sp})$ and $f(\text{sp})$ (Supplementary Material 2) have been increased by a factor
407 of 4, consistent with the modifications adopted for the experimental samples described above.

408 By default in THERMOCALC, uncertainties on the interaction energies are set to $\pm 2 \text{ kJ mol}^{-1}$ (Powell
409 and Holland 2008). For the $a-x$ models used here, uncertainties on many interaction energies are not
410 well constrained and therefore we chose initially to retain this approximation during the assessment of
411 avP calculations on the experimental database. Once the final set of modified $a-x$ models had been
412 selected, we noted that avP calculations produced low σ_{fit} ($\ll 1.0$) for many experimental runs and
413 natural samples, implying that the overall input uncertainties were in many cases excessive. We
414 therefore reduced the uncertainties on interaction energies to $\pm 1 \text{ kJ mol}^{-1}$, which increased to σ_{fit} values
415 generally closer to 1.0 (Supplementary Material 1 and Table 3) and decreased $\sigma_{\bar{P}}$ by about 0.2 kbar.
416 Such generalised uncertainties on interaction energies were adopted as they represent the best
417 compromise to simultaneously obtain low $\sigma_{\bar{P}}$ and successful prediction of experimental pressures. In
418 fact, this might be a realistic estimate of the uncertainties on the most critical parameters, but many of
419 the less influential parameters are much less well known.

420 During avP calculations applied to experimental samples, the values of P_0 and T_0 were set to T_{exp}
421 and P_{exp} (we note that uncertainties in the measured pressure and temperature of the experimental runs
422 can be as high as $\pm 0.5 \text{ kbar}$ and $\pm 20^\circ\text{C}$). In the case of natural samples, P_0 was set to 4 kbar,
423 representing the middle of the range of crustal pressure. For both experimental and natural samples, the
424 window within which the avP algorithm should search for the P_r values was set to $P_0 \pm 20 \text{ kbar}$.
425 Methods for choosing T_0 for natural samples are discussed in a later section.

426 **AvP calculations on the experimental database**

427 In this section we report the results of avP calculations on the experimental database using the refined
428 thermodynamic models and the associated uncertainties adopted in this work.

429 **AvP output: \bar{P} , $\sigma_{\bar{P}}$ and σ_{fit}**

430 **Spinel + clinopyroxene + olivine + plagioclase equilibria (SCOIP).** The results of calculations using
431 SCOIP on the 62 experimental samples are good, as shown in Figure 5a (see also Supplementary
432 Material 1). 65% of the calculations passed the χ^2 test (i.e., $\sigma_{\text{fit}} < 1.49$), which indicates that the
433 experimental mineral compositions and the thermodynamic data, with their associated uncertainties,
434 can be safely combined in the averaging procedure. Slightly higher σ_{fit} values (no more than 20%
435 higher than the cut-off imposed by the χ^2 test) might still be considered acceptable, especially since $\sigma_{\bar{P}}$
436 is increased by a factor equal to σ_{fit} when the χ^2 test is failed. This would increase the proportion of
437 acceptable calculations to 82%. For the calculations that passed the χ^2 test, $\sigma_{\bar{P}}$ varies between 0.80 and
438 2.10 kbar, with a mean value of 1.24 kbar. Considering all the calculations, the experimental pressures
439 are predicted within $\sigma_{\bar{P}}$ for 67% of the samples, consistent with the expectation that the interval $\bar{P} \pm \sigma_{\bar{P}}$
440 will encompass the experimental pressure 68% of the time. The accuracy of avP can also be
441 demonstrated by calculating SEE for \bar{P} . For the selected experimental dataset the calculated \bar{P} yields a
442 SEE of ± 1.37 kbar (± 1.20 kbar considering only calculations with χ^2 test passed), which is small
443 relative to the typical values obtained by the existing geobarometers for gabbroic assemblages. For
444 example, the clinopyroxene-orthopyroxene barometer of Putirka (2008) (his eq. 39) has SEE = ± 2.8
445 kbar using his experimental database and SEE = ± 2.7 kbar using our filtered experimental database
446 (see Supplementary Material 1) in the P range 0.5–10 kbar.

447 No dependence of the pressure residuals ΔP ($\bar{P} - P_{\text{exp}}$) on P_{exp} (Fig. 5a) or T_{exp} (Fig. 6a) is observed,
448 suggesting that the P - and T -dependency of the adopted a - x relations is appropriate for the purpose of
449 avP calculations. Furthermore, the residuals are not dependent on the composition of plagioclase (An =
450 25–94 mol %), olivine (Fo = 61–95 mol %) or spinel ($\text{Fe}^{3+}/[\text{Fe}^{3+} + \text{Al} + \text{Cr}] = 0.00$ – 0.90 ; Fig. 6b). For
451 clinopyroxene, the residuals do not appear to vary with composition, although a slight dependence of

452 $\sigma_{\bar{P}}$ with mg# might be inferred (Fig. 6c), given the apparent higher values of $\sigma_{\bar{P}}$ when mg# < 0.80. No
453 correlations or systematic deviations in ΔP are observed between hydrous and nominally-anhydrous
454 experiments (Supplementary Material 1). Overall, the application of avP to SCOIP with the adopted
455 thermodynamic model produces satisfactory results for a wide range of experimental compositions, and
456 can therefore be applied with confidence to natural mafic and ultramafic rocks that contain spinel,
457 clinopyroxene, olivine and plagioclase. More accurate $\text{Fe}^{3+}/\Sigma\text{Fe}$ estimations for clinopyroxene and
458 spinel would likely increase the accuracy of the calculated pressures. However, the good avP results
459 obtained in our calculations suggest that the errors related to Fe^{3+} calculated from stoichiometry do not
460 generally have a significant deleterious effect on \bar{P} .

461 **Clinopyroxene + olivine + plagioclase equilibria (COIP).** Fig. 5b shows the results of avP
462 calculations with COIP equilibria on the same experimental dataset used for SCOIP. Most of the
463 calculations passed the χ^2 test ($\sigma_{\text{fit}} < 1.73$) and $\sigma_{\bar{P}}$ varies between 0.93 kbar and 3.20 kbar. Results for
464 samples at $P_{\text{exp}} \geq 5$ kbar show a good agreement with the experimental pressures and low calculated
465 uncertainties (Fig. 5b). However, pressures are overall predicted within the uncertainties for a smaller
466 proportion of experiments (54%), which is linked to a systematic overestimation of ~ 1.5 kbar for
467 samples equilibrated at $P_{\text{exp}} < 5$ kbar. Comparable results were obtained on the entire dataset of 209
468 experiments containing COIP (Supplementary Material 6). We suspect that the source of the
469 overestimation at lower pressures could be an excess amount of Al^{IV} in clinopyroxene related to growth
470 rate (Mollo et al. 2010), which is expected to have a relatively larger effect where equilibrium values
471 for Al^{IV} are low, i.e., at low pressures. This issue does not significantly influence the accuracy of \bar{P}
472 when reactions involving spinel end-members are included in the independent set (i.e., SCOIP).

473 **AvP output: residuals and correlations**

474 In this section we examine the results of the least-squares minimization for the individual end-members
475 involved in the reactions and the correlation between the six SCOIP reactions. The low σ_{fit} values
476 obtained for most of the calculations on the experimental database indicate that, in general, the
477 calculated end-member activities and enthalpies do not have to be adjusted outside their uncertainties to
478 obtain a consistent value of P (i.e. \bar{P}) for the six reactions. In the avP diagnostics this is shown by the
479 low uncertainty-normalised residuals in log (activity) (e_i^*) and the enthalpy (H_i^*) for each end-member i
480 (see Supplementary Material 2). For calculations that passed the χ^2 test, $|e_i^*|$ is always <2.5 and $|H_i^*|$ is
481 always <1.0 for all end-members. No systematic outliers appear for the calculations on different
482 samples, with both e_i^* and H_i^* scattering randomly around zero. The only exception is for the end-
483 member acmite, which shows e_i^* varying between -0.1 and 2.5 and H_i^* between -0.2 and 1.4 . Possibly,
484 such systematic positive residuals indicate a slight systematic overestimation of Fe^{3+} in clinopyroxene
485 as determined by stoichiometry. For calculations that did not pass the χ^2 test no end-member is
486 identifiable that systematically shows high residuals (i.e., $|e_i^*| > 2.5$ and $|H_i^*| > 1.0$). The diagnostic h_i
487 (see Supplementary Material 2) shows that for most of the experimental samples no single end-member
488 has a prevailing influence on the calculated \bar{P} . The end-members that are most likely to be influential
489 are cats, cenh, sp, herc and jd. However, in a typical calculation such as that illustrated in the
490 Supplementary Material 2, if the uncertainty on $\ln a_i$ of any single end-member i is doubled, the value
491 of \bar{P} changes by ≤ 0.2 kbar, except in the case of cats for which it is reduced by 0.4 kbar.

492 Supplementary Material 7 shows the correlation matrixes of the enthalpies of reactions (ΔH_r), $\ln K_r$ and
493 P_r , for the calculations on the experimental database. The matrix of ΔH_r values is the same in each
494 calculation. ΔH_r values for the reaction pairs r1-r4 (0.729), r2-r4 (0.746) and r3-r5 (0.741) are
495 particularly highly correlated due to their shared end-members. These reaction pairs also show
496 relatively highly correlated values of $\ln K_r$, but the magnitude of the correlations in $\ln K_r$, and
497 consequently in P_r , are strongly dependent on phase compositions. In a calculation for which P_{r3} and P_{r5}

498 have a correlation coefficient of 0.99, these two reactions amount to a single independent constraint on
499 how the end-member properties should be adjusted to attain a consistent \bar{P} . However, in calculations
500 involving other phase compositions these reactions show only weak to moderate correlation.

501 **Sensitivity of AvP to P_0 and T_0**

502 Calculations with avP require an initial estimate of pressure (P_0) and temperature (T_0) to be specified.
503 These are needed for the initial calculation of end-member enthalpies and activities. We performed a
504 test to investigate how different assumptions about P_0 and T_0 may influence the calculated SCOIP
505 equilibria and the subsequent results of avP. Multiple calculations with avP on the selected experiments
506 have been performed, firstly varying P_0 in the range 0–15 kbar (with $T_0 = T_{\text{exp}}$) and then varying T_0 in
507 the range 600–1400°C (with $P_0 = P_{\text{exp}}$). The first tests showed a negligible effect of P_0 on the calculated
508 pressure. For the sample ShR658 ($P_{\text{exp}} = 2.0$ kbar, $T_{\text{exp}} = 1075$ °C; Supplementary Material 1), for
509 example, a substantial difference of 15 kbar in P_0 produces a difference in the calculated \bar{P} of 0.01 kbar
510 ($\bar{P} = 2.94 \pm 1.62$ kbar if $P_0 = 0$ kbar; $\bar{P} = 2.95 \pm 1.62$ kbar if $P_0 = 15$ kbar), with no difference in the
511 quality of the fit. Similar results have been observed for other experimental samples in the dataset.

512 On the other hand, different assumptions about T_0 variably affect the calculated \bar{P} . An example is
513 reported in Fig. 7, which shows the results for the same experiment ShR658. In this case, the
514 temperature dependence of \bar{P} , $d\bar{P}/dT_0$, corresponds to -0.012 kbar/°C. More generally the test on the
515 experimental dataset showed that $d\bar{P}/dT_0$ is always negative. Furthermore, the absolute value of the
516 dependence of \bar{P} on T_0 is smaller for greater values of \bar{P} , with $d\bar{P}/dT_0$ varying from -0.016 to -0.002
517 kbar/°C (Fig. 8). This indicates that, for natural assemblages equilibrated at high pressures, even errors
518 related to crude temperature estimates (e.g., from geological/petrological inferences) will not
519 significantly increase the uncertainty on the calculated pressure (i.e., additional ± 0.4 – 0.6 kbar for
520 errors on T_0 on the order ± 150 – 200 °C).

521 For assemblages equilibrated at relatively low pressures (<5kbar), uncertainties on T_0 on the order of
522 150–200°C might increase the total pressure uncertainties by 1.5–2.5 kbar, which indicates the need of
523 reliable methods to estimate equilibration temperature when applying avP to natural samples. One such
524 a method is to carry out the avP procedure at multiple temperatures, and observe the variation in the
525 value of σ_{fit} . If the variation of σ_{fit} with T_0 follows a polynomial function of the type $\sigma_{\text{fit}} = a'T_0^2 + b'T_0$
526 $+ c'$ (where a' , b' and c' are constants and $a' > 0$, $b' < 0$), and a minimum of σ_{fit} is clearly identifiable
527 (e.g., Fig. 7), then the value of T_0 that gives the minimum value of σ_{fit} can be considered to give better
528 estimates of T . Figure 7 shows a successful example of such a method, in which the estimated
529 temperature ($T_{\sigma_{\text{fit}}}$; 1050°C) is close to the measured experimental temperature, T_{exp} (1075°C). We
530 applied such a method to the selected experimental dataset, firstly to evaluate the reliability of the
531 estimated temperatures and secondly to assess how the possible errors associated with these estimates
532 affect the calculated \bar{P} . For some experiments the variation of σ_{fit} with T_0 is such that a minimum in σ_{fit}
533 cannot be identified and therefore these experiments ($n = 25$) have been removed from the test dataset.
534 The results of the test are reported in Figure 9, which shows that $T_{\sigma_{\text{fit}}}$ predicts the experimental
535 temperatures with a SEE of $\pm 93^\circ\text{C}$ ($n = 37$). Application of avP to the same experimental dataset,
536 adopting $T_0 = T_{\sigma_{\text{fit}}}$ instead of $T_0 = T_{\text{exp}}$, still produces acceptable calculated pressures (cf. Figs. 9b and
537 d).

538 We note that $\sigma_{\bar{P}}$ also varies with T_0 , and the minima in $\sigma_{\bar{P}}$ and σ_{fit} do not generally occur at the
539 same T_0 . This suggests that temperature is not strongly constrained in the sample. We find empirically
540 that the optimal T_0 is usually better predicted by the minimum in σ_{fit} , even when σ_{fit} passes the χ^2 test
541 over a wide range of temperatures. The use of $T_{\sigma_{\text{fit}}}$ as a possible T estimate can be justified in that σ_{fit}
542 characterises the compatibility of the different independent reactions, and T_0 contributes to the
543 formulation of those reactions without being formally included in the covariance matrix summarised by
544 $\sigma_{\bar{P}}$.

545 In conclusion, P_0 does not have significant effects on the calculated \bar{P} at any condition and therefore
546 we suggest the initial estimate to be always set to 4 kbar (i.e., mid-crustal values). On the other hand,
547 the effect of T_0 might be in some cases significant, but it can be easily evaluated for each sample by
548 running multiple avP calculations at different T_0 . If no other thermometer can be applied that gives T
549 uncertainties < 90 °C, then $T_0 = T_{\sigma_{\text{fit}}}$ can be used to produce the best estimate of P.

550 **Application to natural rocks**

551 Given the success of tests on the experimental dataset, we next applied the refined thermodynamic
552 models to estimating the pressure of formation of igneous rocks, choosing samples containing spinel,
553 clinopyroxene, olivine and plagioclase. We carefully selected two groups of mafic/ultramafic xenoliths
554 with texture, mineralogy and mineral chemistry suggesting good attainment of chemical equilibrium.
555 Compositions of minerals in these natural samples are reported in Fig. 2. The first group of xenoliths,
556 from the island of Dominica, Lesser Antilles, provides an example of the application of avP to gabbroic
557 rocks from a mid-upper island arc crust. The second group, which was described and analysed by
558 McGuire (1994) and McGuire and Mukasa (1997), is from the Wikieup volcanic field, Arizona, and
559 provides an example of gabbroic and ultramafic xenoliths from the lower continental crust. Unless
560 otherwise specified, the following discussion refers to calculations with SCOIP.

561 **Gabbroic xenoliths from Dominica, Lesser Antilles**

562 **Background, sample description and analyses.** The Lesser Antilles is an active volcanic arc resulting
563 from westwards subduction of the South American plate beneath the Caribbean plate (Macdonald et al.
564 2000). The volcanic products, which vary from basalts to andesites and dacites, host a large number of
565 plutonic xenoliths showing a variety of mineralogical compositions (Arculus and Wills 1980; Tollan et
566 al. 2012; Stamper et al. 2014b; Cooper et al. 2016). The ubiquity of these xenoliths provides an
567 exceptional opportunity to investigate the nature of the crust and the evolution of magmatic systems in

568 an island arc setting, but such an investigation requires precise estimates of the pressure of formation of
569 the xenoliths. Given that the estimated thickness of the crust in the Lesser Antilles varies between 25
570 and 35 km (cf., Boynton et al. 1979; Christeson et al. 2008; Kopp et al. 2011), the gabbroic xenoliths
571 are likely to have equilibrated at pressures between ~1 and ~8 kbar. Therefore, only geobarometric
572 estimates with low associated uncertainties (<2 kbar) can provide useful insight into the vertical
573 distribution of mafic lithologies in the crust.

574 Dominica, in the centre of the arc, is mainly composed of Pliocene to recent volcanic rocks (Lindsay
575 et al. 2005; Howe et al. 2015). Plutonic xenoliths, mostly consisting of gabbroic assemblages with or
576 without olivine, orthopyroxene and hornblende, have been found as clasts in young, reworked
577 sediments from the southern part of the island. For the purpose of testing the applicability of the avP
578 method, we selected three xenoliths containing the assemblage SCOIP, which have been collected from
579 marine deposits close to the village of Petite Savanne (southern Dominica), on the edge of the
580 Foundland volcanic centre. A schematic description of these samples, highlighting the main textural
581 and mineralogical features, is shown in Supplementary Material 8. Note that xenolith DC93 is
582 composed of two parts (a coarse-grained troctolite, DC93a, attached to a heterogranular Ol-
583 gabbroite, DC93b).

584 Details of electron microprobe analyses and results are reported in the Supplementary Material 5.
585 Compositions of olivine, spinel, clinopyroxene and plagioclase (Fig. 2) are similar to those found in
586 other cumulate xenoliths from the Lesser Antilles (e.g., Tollan et al. 2012; Stamper et al. 2014b;
587 Cooper et al. 2016). Coexisting olivine and plagioclase show relatively low forsterite contents (68–73
588 mol% Fo) and high anorthite contents (87–93 mol% An) respectively; clinopyroxenes are augites
589 showing a limited range of mg# (78–82 %); spinels are titanomagnetites with TiO₂ contents of 7.6 to
590 9.2 wt%. Grain-to-grain and core-rim variations for olivine, spinel, pyroxenes and plagioclase are
591 small, suggestive of well-equilibrated assemblages. The largest compositional variations are observed

592 in samples DC91 and DC93b (e.g., $\text{Al}_2\text{O}_3^{\text{cpx}} = 1.9\text{--}3.7$ wt% for DC91). Clinopyroxene in DC93b shows
593 a poikilitic texture with anhedral spinel inclusions and a slight patchy zoning mainly limited to Al^{IV}
594 ($0.07\text{--}0.13$ apfu) and Fe^{3+} ($0.03\text{--}0.04$ apfu). For the purpose of thermobarometry, only the
595 compositions of adjacent crystal rims have been used (Fig. 10).

596 **Average P calculations.** Mineral compositions reported in Supplementary Material 5 have been used
597 to calculate \bar{P} , adopting the thermodynamic models and calculation methods described above (see also
598 Supplementary Material 2 for the example of an avP calculation on sample DC90). T_0 has been
599 estimated by running a calculation every 10°C in the range $700\text{--}1200^\circ\text{C}$, and selecting the temperature
600 that produce the lowest σ_{fit} ($T\sigma_{\text{fit}}$) (see above and Fig. 7). Temperatures calculated with this method
601 ($870\text{--}1000^\circ\text{C}$) and results of avP calculations are reported in Table 3. All calculations produced low σ_{fit}
602 values ($0.81\text{--}1.54$) and $\sigma_{\bar{P}}$'s for each sample are low ($0.91\text{--}1.55$ kbar). The lowest $\sigma_{\bar{P}}$ is obtained for
603 the ol-gabbro-norite DC90 ($\bar{P} = 3.58 \pm 0.91$ kbar), which is the sample with the smallest compositional
604 heterogeneities and best-equilibrated texture (e.g. $\sim 120^\circ$ grain boundaries). The cpx-amph-gabbro-norite
605 DC91 contains also minor orthopyroxene, which allowed for application of cpx-opx thermometry
606 (Putirka 2008; Table 3). T estimated with this method is in agreement with $T\sigma_{\text{fit}}$ (896 vs 910°C),
607 although the observed $K_{\text{D}}(\text{Fe-Mg})^{\text{cpx-opx}}$ is below the range considered by Putirka (2008) to demonstrate
608 equilibrium (i.e., 0.79 vs. 1.09 ± 0.14). Pressure estimated for this sample (at $T_0 = T\sigma_{\text{fit}}$) is 1.94 ± 1.09
609 kbar. For all four Dominica samples, uncertainty-normalised residuals are low ($|e_i^*| < 1.5$, $|H_i^*| < 0.27$),
610 with acmite showing somewhat higher values ($|e_{\text{acm}}^*| < 2.1$, $|H_{\text{acm}}^*| < 0.87$), as observed for the
611 calculations on the experimental database. End-members cats, cenh, and jd show the highest, but still
612 acceptable, h_i values ($h_{\text{cats}} = 0.26\text{--}0.30$, $h_{\text{cenh}} = 0.12\text{--}0.18$, $h_{\text{jd}} = 0.29\text{--}0.35$). Calculations with COIP
613 produced comparable pressure estimates and uncertainties, although \bar{P} is systematically higher (by ≤ 1
614 kbar, Table 3), as previously observed in tests on the experimental database at $P_{\text{exp}} < 5$ kbar.

615 The composite xenolith DC93 contains two assemblages (DC93a, DC93b; Supplementary Material
616 8 and Fig. 10) that show different textures and degrees of chemical and textural equilibrium despite
617 likely forming at the same pressure. Ol-gabbro-norite DC93b show somewhat larger compositional
618 heterogeneities and lower Al^{IV} and Fe^{3+} in clinopyroxene compared to troctolite DC93a
619 (Supplementary Material 5), which is most likely related to a difference in bulk composition. This
620 provides an opportunity to test the effect of such heterogeneities on avP calculations with SCOIP.
621 Calculated pressures for samples DC93a and DC93b are in good agreement (2.67 ± 1.35 kbar vs, 2.81
622 ± 1.55 kbar), supporting the robustness of avP.

623 The geobarometric results obtained for these samples give an estimate of the pressures of magma
624 storage and crystallization in southern Dominica. Overall, the calculated pressures are in agreement
625 with the range inferred by observational and experimental studies on magmas and cumulate rocks from
626 the Lesser Antilles (i.e. mostly 1.0–5.0 kbar for gabbroic assemblages; c.f., Pichavant and Macdonald
627 2007; Stamper et al. 2014b; Melekhova et al. 2015; Cooper et al. 2016). Pressure calculated for the
628 cpx-amph-gabbro-norite DC91 (1.94 ± 1.09 kbar) agree very well with the pressures of orthopyroxene
629 and plagioclase crystallization in lavas from southern Dominica (~ 2 kbar; Gurenko et al. 2006),
630 estimated using melt inclusion data and H_2O-CO_2 solubility models. Considering an average density
631 ($\bar{\rho}$) of the crust of 2800 kg/m^3 , the calculated $\bar{P} \pm \sigma_{\bar{P}}$ for this cpx-amph-gabbro-norite corresponds to a
632 depth of 7.1 ± 4.0 km, which agrees with the estimated hypocentres of deep volcanic earthquakes in the
633 nearby Plat Pays volcanic complex (2–6 km; Lindsay et al. 2003), thought to result from pressurization
634 of the magma reservoirs. Note that application of the cpx-opx geobarometer of Putirka (2008; his eq.
635 39) to the same sample DC91, which produces an estimated value of 1.6 kbar (as stated above, $K_D[Fe-$
636 $Mg]^{cpx-opx} < 0.95$ potentially flags incomplete equilibration), is accompanied by a model uncertainty of
637 ± 2.8 kbar (Putirka 2008). This corresponds to an uncertainty of ± 10.2 km in the estimated depth of
638 formation.

639 Pressures estimated for ol-gabbro norite DC90 and troctolite DC93a correspond to depths of $13.0 \pm$
640 3.3 km and 9.7 ± 4.9 km respectively. Based on the seismic reflectors identified for the crust north of
641 Dominica (Kopp et al. 2011), such depths correspond to the middle-to-upper crust transition. Figure 11
642 shows how the thermobarometric results on all three xenoliths compare with the geophysical model of
643 Kopp *et al.* (2011). Note that among the existing geobarometers calibrated for basaltic compositions,
644 only the single-clinopyroxene formulations of Nimis (1995) and Putirka (2008) could be applied to
645 these assemblages². Application of equation 32a of Putirka (2008) produces $P = 9.4 \pm 3.1$ kbar for the
646 sample DC90 and $P = 9.9 \pm 3.1$ kbar for the sample DC93a, which corresponds to depths of 34.2 ± 11.3
647 km and 36.0 ± 11.3 km, respectively, exceeding the crustal thickness in this area (27 km; Kopp et al.
648 2011). Overall, the successful application of avP to these samples suggest that future studies on a more
649 extensive xenolith suite from Dominica and elsewhere in the Lesser Antilles arc, combined with
650 petrological and geophysical observations, may help to constrain the lithological structure of the crust
651 with a precision not previously achievable.

652 **Mafic and ultramafic xenoliths from Wikieup, Arizona**

653 **Background and sample selection.** The Wikieup volcanic field is located in the Arizona Transition
654 Zone (south-western United States), between the Basin and Range Province (BRP), a region of thinned
655 continental crust (20–30 km), and the Colorado Plateau (CP), characterized by a thicker crust (40-50
656 km) that remained tectonically stable over the Phanerozoic (e.g., Bashir et al. 2011). The Wikieup
657 volcanics host a large number of mafic and ultramafic xenoliths thought to represent portions of the
658 underlying deep crust and upper mantle. Some of the xenoliths analysed by McGuire (1994) and later
659 by McGuire and Mukasa (1997) contain the assemblage SCOIP, making them suitable for application
660 of the avP method. Quantifying the thickness and nature of the crust in the Arizona Transition Zone is
661 important to interpret the tectonic and magmatic events that shaped the lithosphere of the southwestern

² Note that there are no tests for equilibrium (e.g., Putirka 2008) for clinopyroxene-only models.

662 US. Consequently a large number of geophysical investigations have been carried out in this area (e.g.,
663 Hauser et al. 1987; Parsons et al. 1992; Bashir et al. 2011 and reference therein), providing a context
664 for our thermobarometric results. McGuire (1994) presented such an approach using gabbroic
665 xenoliths, but due to the lack of reliable geobarometers for such assemblages, the estimated pressures
666 were poorly constrained.

667 Of the xenoliths described and analysed by McGuire (1994) and McGuire and Mukasa (1997) we
668 used four samples, from Tertiary alkaline olivine-basalts (Wilshire 1990) on the eastern flank of the
669 Hualuapai Mountains, that contain the assemblage SCOIP and for which mineral compositions are
670 available. Samples WK2-9 and WKa-49a are ol-gabbros; samples WK1-24 and WK1-63 are
671 plagioclase-peridotites. The latter show higher plagioclase and lower olivine contents with respect to
672 typical mantle peridotites, but show compositional and textural features that distinguish them from the
673 ol-gabbros (McGuire and Mukasa 1997). Both pl-peridotites and gabbros show allotriomorphic-
674 granular textures but pl-peridotites are also often foliated. Some textural evidence of early-stage
675 reactions between plagioclase and olivine grains has been reported (McGuire and Mukasa 1997), but
676 compositional heterogeneity is limited to some plagioclase grains that show slight normal zoning (<
677 2% An; McGuire and Mukasa 1997). Phase compositions are shown in Fig. 2 and detailed microprobe
678 analyses are reported in McGuire (1994) and McGuire and Mukasa (1997). Olivine shows a limited
679 variation in composition (Fo₇₇₋₇₉); spinel is Al-rich ($Al/[Al+Cr+Fe^{3+}] = 0.90-0.97$) with low TiO₂
680 contents (0.11–0.45 wt%); clinopyroxenes are augites with Al₂O₃ contents varying from 6.91 to 7.49
681 wt%; and plagioclase varies between An₅₂₋₅₃ (pl-peridotites) and An₅₈₋₇₁ (ol-gabbros).

682 **AvP calculations.** All avP calculations passed the χ^2 test and the calculated uncertainties are notably
683 low (0.95–1.00 kbar; Table 3). Estimated pressures are similar within the uncertainties for all four
684 samples, but ol-gabbros shows some apparent lower values than pl-peridotites (Table 3). Similar
685 estimates are obtained using COIP, with no significant difference in the calculated uncertainties (Table

686 3), as previously observed at $P > 5$ kbar for our experimental test (cf., Figs. 5a and b). Note that ol-
687 gabbro WK2-9 and pl-peridotite WK1-63 also contain orthopyroxene, allowing $T_{\sigma_{\text{fit}}}$ to be compared
688 with the results of cpx-opx thermometry (Putirka, 2008). Observed $K_D(\text{Fe-Mg})^{\text{cpx-opx}}$ (0.94–0.95) are
689 within the range recommended for the application of such thermometer (Putirka, 2008). Temperatures
690 for the ol-gabbro WK2-9 agree fairly well (940°C vs. 887°C), while temperatures for the pl-peridotite
691 WK1-63 disagree by about 200°C (Table 3). However, calculations for WK1-63 using T_0 estimated
692 from cpx-opx thermometry ($T_0 = 953^\circ\text{C}$ instead of $T_0 = 1150^\circ\text{C}$) produce relatively small differences in
693 the calculated pressure (7.78 ± 0.86 kbar vs. 6.88 ± 0.97 kbar).

694 When calculated pressures for all four samples are converted to depths (Table 3), the resulting
695 values are consistent with those suggested by McGuire (1994) and McGuire and Mukasa (1997) based
696 on: i) the observation that many composite xenoliths show contact between gabbros and mantle
697 peridotites, suggesting an origin from the crust-mantle boundary (McGuire 1994) and ii) projection of
698 calculated temperatures onto the regional geotherm (100 mW/m²; McGuire and Mukasa 1997). Our
699 calculations place more robust petrological constraints on the equilibration depths for comparison with
700 recent geophysical investigations (Bashir et al. 2011). Figure 12 shows that the calculated depths for
701 the ol-gabbros and pl-peridotites correspond closely to the lower crust levels in the integrated gravity
702 and seismic model of Bashir *et al.* (2011). Note that the upper-lower crust and crust-mantle boundaries
703 sketched in Fig. 12 should be viewed as transition zones and that the geophysical model has an
704 associated uncertainty, which that is not easily quantifiable (Bashir et al. 2011). The fact that both ol-
705 gabbros and pl-peridotites show comparable depths of formation is in agreement with a model of the
706 lower continental crust characterized by a network of dykes and sills of variably mafic lithologies
707 (McGuire 1994; Parsons et al. 1992), that can be related to the magmatic episodes that brought the
708 xenoliths to the surface. Further application of the avP method to a larger number of xenoliths,

709 combined with detailed modal analyses and integrated geophysical-petrological models, may help to
710 test this hypothesis.

711 **Implications and final recommendations**

712 Application of avP to experimental phase equilibria and well-equilibrated natural rocks has shown, for
713 the first time, that this method can be used to reliably estimate the pressure of formation for mafic and
714 ultramafic igneous rocks. The least squares minimization procedure in avP makes it possible to exploit
715 the power of internally-consistent thermodynamic datasets and mutually compatible $a-x$ relations,
716 taking into account both uncertainties and correlations in end-member properties. It represents an
717 optimal method to deal with the inevitable inaccuracy of thermodynamic models and uncertainties on
718 mineral compositions. We demonstrated that calculations on SCOIP equilibria, using the models
719 refined in this work, yield i) good reproducibility of experimental pressures by avP (P is predicted,
720 within the calculated $\sigma_{\bar{P}}$, for 67% of the experimental database), ii) internal consistency of avP results
721 (i.e., χ^2 test passed) for the majority of phase equilibrium experiments and for the selected well-
722 equilibrated natural assemblages, and iii) low P uncertainties ($\sigma_{\bar{P}}$) as low as 0.9 kbar) for the same
723 assemblages. Pressure estimates with such low uncertainties can help to characterise the key
724 petrological processes and trace the vertical distribution of mafic and ultramafic lithologies in the upper
725 30–40 km of the lithosphere. Integration with geophysical data for the same areas affords potentially
726 powerful insights into crustal structure.

727 Modifications of the existing $a-x$ models for clinopyroxene (Green et al. 2016) and spinel (Jennings
728 and Holland 2015) were necessary to avoid some observed pressure deviations resulting from initial
729 calculations with SCOIP. At present we believe that the thermodynamic models used in this work are
730 applicable to any gabbroic rock, pyroxenite or peridotite that contains olivine + clinopyroxene +
731 plagioclase in its equilibrium assemblage, regardless of geodynamic setting. This is supported by i) the
732 robustness of a multiple-reaction approach based on an internally consistent thermodynamic dataset

733 (Holland and Powell 2011) and ii) the successful tests on a database of phase equilibrium experiments
734 that cover a wide range of basaltic and peridotitic compositions (e.g., Figs. 5 and 6), both hydrated and
735 nominally-anhydrous. We believe the method can be applied also to volcanic rocks, for which,
736 however, the selection of equilibrium phase compositions might not be straightforward. Future test
737 cases, in which a detailed textural and chemical characterization of the phenocryst assemblages will be
738 coupled with the application of avP and its fit diagnostics, will help to test this hypothesis.

739 For equilibrated assemblages including spinel, calculations with SCOIP equilibria are always
740 preferred, since COIP slightly overestimate pressures (by 1–2 kbar) at upper-middle crust conditions (P
741 < 5 kbar). Calculations with SCOIP should be possible for a large number of mafic/ultramafic rocks, as
742 spinel is a ubiquitous phase in many natural gabbroic and basaltic rocks, and sometimes coexists with
743 plagioclase in shallow mantle peridotites/pyroxenites. Calculations with COIP, on the other hand, are
744 more likely to be needed for plagioclase-only peridotites/pyroxenites. Importantly, the latter typically
745 form at lower crust and uppermost mantle conditions ($P > 5$ kbar; e.g., Kelemen et al. 2007; Borghini et
746 al. 2011), corresponding to those pressures at which COIP calculations show low uncertainties and are
747 consistent with SCOIP based on tests on experimental (Fig. 5) and natural samples (Table 3).

748 The reliability of avP calculations on natural samples relies on the attainment of equilibrium in the
749 assemblage. If the assemblage shows slight intra- or inter-grain chemical heterogeneities, care should
750 be taken to select mineral compositions that are likely to be at equilibrium, as shown by the application
751 to the cumulate xenoliths from Dominica (Fig. 10). Secondly, when using our sets of independent
752 reactions (Fig. 4), temperature estimates for the sample need to be well constrained if the assemblage
753 equilibrated at $P < 5$ kbar. In order to estimate the true T of equilibration, it may be possible to use
754 existing thermometers, such as the cpx-opx thermometer of Putirka (2008) or the Ca-in-olivine
755 thermometer of Shejwalkar and Coogan (2013). Otherwise we suggest running avP calculations at

756 different temperature T_0 , and select that which produces the lowest residuals (σ_{fit}) in the avP
757 minimization routine (Fig. 7).

758 Further work is planned to extend the avP approach to a wider range of phase assemblages. This will
759 entail testing and potential modification of thermodynamic models for orthopyroxene, amphibole and
760 ilmenite, which are also common constituents of mafic and ultramafic igneous rocks in various settings.
761 Using such a set of models, it will be possible to obtain mutually consistent pressure estimates for a
762 wide range of igneous rocks.

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- 1062

1063

Figure captions

1064

1065 **Figure 1.** *P-T* conditions of phase equilibrium experiments in basaltic and peridotitic systems
1066 that contain at least olivine, clinopyroxene and plagioclase in their phase assemblages. Source
1067 references for this database are reported in Supplementary Material 1.

1068

1069 **Figure 2.** Composition of the mineral phases used for avP calculations on experimental and
1070 natural samples. (a) Pyroxene quadrilateral projection (Morimoto 1988). (b) Forsterite content in
1071 olivine plotted vs anorthite content in plagioclase. (c) Trivalent cation plot for spinels [$\text{Fe}^{3+} =$
1072 $\text{Fe}^{3+}/(\text{Fe}^{3+}+\text{Al}+\text{Cr})$, $\text{Al} = \text{Al}/(\text{Fe}^{3+}+\text{Al}+\text{Cr})$, $\text{Cr} = \text{Cr}/(\text{Fe}^{3+}+\text{Al}+\text{Cr})$].

1073

1074 **Figure 3.** Variation of Al^{IV} in clinopyroxene with pressure in the experimental database. Note the
1075 wide range of Al^{IV} for the experiments at 0.001 kbar, which is larger than the variations for the
1076 entire dataset at higher pressures.

1077

1078 **Figure 4.** Independent set of reactions and avP results at $T_0 = 870$ °C for the spinel +
1079 clinopyroxene + olivine + plagioclase (SCOIP) assemblage in olivine-gabbro DC90
1080 (Dominica, Lesser Antilles). The reactions are linearised at $T_0 = 870$ °C and $P_0 = 4$ kbar. In the
1081 minimization procedure, avP makes minimal adjustments to the enthalpies and activities of end-
1082 members, given their covariance matrices, in order for the reactions to coincide at \bar{P} . Reactions
1083 derived from the initial input data, i.e. the dataset enthalpies and the activities at the measured
1084 compositions, are shown as solid curves. Reactions after minimization are shown as dashed curves.
1085 The box at the bottom shows the pressures initially obtained from each individual reaction with their
1086 associated uncertainties, at $T_0 = 870$ °C. For end-member abbreviations see Table 2. A detailed step-
1087 by-step demonstration of the same calculation is reported in the Supplementary Material 2.

1088

1089

1090 **Figure 5.** Results of avP calculations on the refined database of experiments containing the
1091 subset SCOIP in their run products. Calculations have been performed with either **(a)** SCOIP
1092 equilibria (i.e., 6 reactions) or **(b)** COIP equilibria (3 reactions), using the thermodynamic model
1093 refined in this work. Error bars are $\sigma_{\bar{P}}$ (see text). The labels on the x-axes are the experiment
1094 index numbers, ordered as in the Supplementary Material 1 (i.e., increasing P_{exp})

1095

1096 **Figure 6.** Residuals in calculated \bar{P} relative to P_{exp} , for SCOIP equilibria, plotted vs. **(a)** T_{exp} , **(b)**
1097 spinel $\text{Fe}^{3+}/[\text{Fe}^{3+}+\text{Al}+\text{Cr}]$ and **(c)** clinopyroxene mg#. Dataset and symbols as in Fig. 5.

1098

1099 **Figure 7.** Results of multiple avP calculations at different temperatures T_0 for the experimental
1100 sample ShR658 (Husen et al. 2016). $P_0 = P_{\text{exp}}$ for all the calculations. The temperature
1101 dependence of the calculated pressure ($d\bar{P}/dT_0$) for this sample is indicated in the top box. $T\sigma_{\text{fit}}$
1102 is the value of T_0 that produces the avP result with the lowest σ_{fit} ($\bar{P}\sigma_{\text{fit}}$). Note, with reference to
1103 the uncertainty envelope (grey envelope in the top box), that outside the T range in which the χ^2
1104 test is passed, the uncertainty is multiplied by σ_{fit} .

1105

1106 **Figure 8.** Temperature dependence of avP ($d\bar{P}/dT_0$) plotted vs. calculated pressures (at $T_0 = T_{\text{exp}}$),
1107 for the selected experimental dataset used in this work. The higher the calculated value of \bar{P} , the
1108 less sensitive this value is to the choice of T_0 .

1109

1110 **Figure 9.** Temperature estimates derived by minimising σ_{fit} ($T\sigma_{\text{fit}}$) for the experimental dataset
1111 and their effect on the calculated \bar{P} . **(a)** Comparison between $T\sigma_{\text{fit}}$ and the true temperature T_{exp} .
1112 **(b,c)** Test of $\bar{P} \pm \sigma_{\bar{P}}$ calculated using either **(b)** $T_0 = T\sigma_{\text{fit}}$ or **(c)** $T_0 = T_{\text{exp}}$. In (c), experiments
1113 shown in black are those for which $T\sigma_{\text{fit}}$ cannot be estimated, as no minimum in σ_{fit} as a function
1114 of T_0 can be identified. In all cases $P_0 = P_{\text{exp}}$. SEE: - Standard Error of Estimate; Pred: -
1115 percentage of calculations that predict P_{exp} within $\sigma_{\bar{P}}$.

1116

1117 **Figure 10.** Photomicrographs (plain polarised light) of the selected xenoliths from Dominica,
1118 with locations of the microprobe analyses used for avP calculations. Results of avP calculations
1119 with SCOLP equilibria are also reported in the boxes at the bottom-left corners.

1120

1121 **Figure 11.** Calculated depths for the three xenoliths from Dominica, Lesser Antilles, compared
1122 with the seismic model of Kopp et al. (2011). Depths calculated as in Table 3. Numbers in the
1123 figure refer to v_p velocities (km/s).

1124

1125 **Figure 12.** Calculated depths (as Table 3) for the ol-gabbro and pl-peridotite xenoliths from
1126 Wikieup, shown in comparison to the density profile across the Arizona Transition Zone (Bashir
1127 et al. 2011). BR – Basin and Range; CP – Colorado Plateau.

1128 **Deposit items**

1129 Supplementary Material 1: Experimental database

1130 Supplementary Material 2: Example of avP calculation and diagnostics

1131 Supplementary Material 3: Thermodynamic dataset and activity-composition models

1132 Supplementary Material 4: Uncertainties on compositional variables for the experimental dataset

1133 Supplementary Material 5: Electron microprobe analyses

1134 Supplementary Material 6: AvP calculations with COIP equilibria on the experimental database

1135 Supplementary Material 7: Correlation matrixes of the reactions

1136 Supplementary Material 8: Description of xenoliths from Dominica, Lesser Antilles

1137

1138 **Table 1.** Abbreviations used in this work

SCOIP	<u>S</u> pinel + <u>C</u> linopyroxene + <u>O</u> livine + <u>P</u> lagioclase, where ‘spinel’ may refer to either magnetite or Cr/Al-spinel.
COIP	<u>C</u> linopyroxene + <u>O</u> livine + <u>P</u> lagioclase
apfu	atoms per formula unit
mg#	Mg/(Mg+Fe ²⁺)
NNO	nickel-nickel oxide buffer
$P_{\text{exp}}, T_{\text{exp}}$	measured pressure, temperature of the experiment
P_0, T_0	initial guesses of pressure and temperature
ΔG_r°	Gibbs free energy difference in kJ for the reaction r among the pure end-members
ν_i	reaction coefficient for end-member i , at P, T
μ_i	total chemical potential for end-member i
μ_i^0	chemical potential of the pure end-member at P, T
R	gas constant (0.0083144 kJ/K·mole)
$W_{(m,n)}^j$	pairwise interaction energy representing non-ideal enthalpy of mixing between end-members m and n in phase j .
a_i	activity of end-member i
K_r	equilibrium constant for reaction r
ΔH_r	enthalpy change for reaction r , at P, T
ΔS_r	entropy change for reaction r , at P, T
ΔV_r	volume change for reaction r , at P, T
P_r	pressure at T_0 for reaction r
\bar{P}	average pressure (at T_0) calculated through the avP algorithm
$\sigma_{\bar{P}}$	calculated 1σ uncertainty for \bar{P}
σ_{fit}	diagnostic value that measures the appropriateness of the average, with σ_{fit}^2 distributed as χ^2 .
e_i^*	uncertainty-normalized residuals in log (activity) for end-member i
H_i^*	uncertainty-normalized residuals in the enthalpy for end-member i
h_i	hat value, i.e., degree of influence of the end-member i on the least-square results
$T_{\sigma_{fit}}$	T_0 that produces the avP result with the minimum value of σ_{fit}
SEE	standard error of the estimate

1139

1140

1141

1142 **Table 2.** Solid solutions and compositional end-members used in this work. Columns SCOIP and
 1143 COIP show which end-members are included in the avP calculations using the SCOIP and COIP
 1144 equilibria, respectively. End-members are described in the Holland and Powell (2011) dataset, using
 1145 the abbreviations shown, with the exception of cenh and cfs.
 1146

Solid solution	End-member	Abbreviation	Formula	SCOIP	COIP
Olivine	forsterite	fo	Mg ₂ SiO ₄	x	x
	fayalite	fa	Fe ₂ SiO ₄	x	x
Spinel	spinel	sp	MgAl ₂ O ₄	x	
	hercynite	herc	FeAl ₂ O ₄	x	
	magnetite	mt	Fe ₃ O ₄	x	
	picrochromite*	picr	MgCr ₂ O ₄		
	ulvöspinel*	usp	Fe ₂ TiO ₄		
Clinopyroxene	diopside	di	CaMgSi ₂ O ₆	x	
	clinoenstatite	cenh	Mg ₂ Si ₂ O ₆	x	x
	clinoferrrosilite	cfs	Fe ₂ Si ₂ O ₆	x	x
	jadeite	jd	NaAlSi ₂ O ₆	x	x
	acmite	acm	NaFeSi ₂ O ₆	x	
	Ca-tschermak	cats	CaAl ₂ SiO ₆	x	x
Plagioclase	high-albite	abh	NaAlSi ₃ O ₈	x	x
	anorthite	an	CaAl ₂ Si ₂ O ₈	x	x
	sanidine*	san	KAlSi ₃ O ₈		

1147 * end-members for which balanced reactions cannot be written. Note that these end-members are
 1148 still considered in the calculations of the activities of the end-members involved in the reactions.
 1149

1150 **Table 3.** Results of avP calculations for the xenoliths from Dominica and Wikieup. For all the
 1151 calculations, $P_0 = 4.0$ kbar and $T_0 = T\sigma_{\text{fit}}$. Uncertainties on compositional variables have been
 1152 propagated from analytical errors, with the corrections adopted in this work (see Supplementary
 1153 Material 5). $T_{\text{cpx-opx}}$ refers to temperature calculated using eq. 37 in Putirka (2008). Depths of the
 1154 xenoliths have been calculated from the estimated pressures, assuming an average density ($\bar{\rho}$) of
 1155 the overlying crust of 2800 g/cm^3 .
 1156

Sample	$T_{\text{cpx-opx}}$ (°C)	$T\sigma_{\text{fit}}$ (°C)	SCOIP equilibria			depth km	COIP equilibria		
			\bar{P} kbar	$\sigma_{\bar{P}}$ kbar	σ_{fit}		\bar{P} kbar	$\sigma_{\bar{P}}$ kbar	σ_{fit}
Dominica:									
DC90	-	870	3.58	0.91	0.81	13.0 ± 3.3	4.21	0.98	0.36
DC91	896	910	1.94	1.09	0.94	7.1 ± 4.0	2.81	1.23	0.40
DC93a	-	980	2.67	1.35	1.48	9.7 ± 4.9	3.66	0.97	0.92
DC93b	-	1000	2.81	1.55	1.54	10.2 ± 5.6	3.76	1.41	1.31
Wikieup:									
WK2-9	887	940	6.53	0.97	0.40	23.8 ± 3.5	6.55	0.99	0.54
WK2-49a	-	1000	6.20	1.00	0.74	22.6 ± 3.6	6.19	1.20	1.17
WK1-24	-	1150	6.54	0.95	0.55	23.8 ± 3.5	6.46	0.99	0.84
WK1-63	953	1150	6.88	0.97	0.48	25.0 ± 3.5	6.86	1.00	0.73

1157

Figure 1

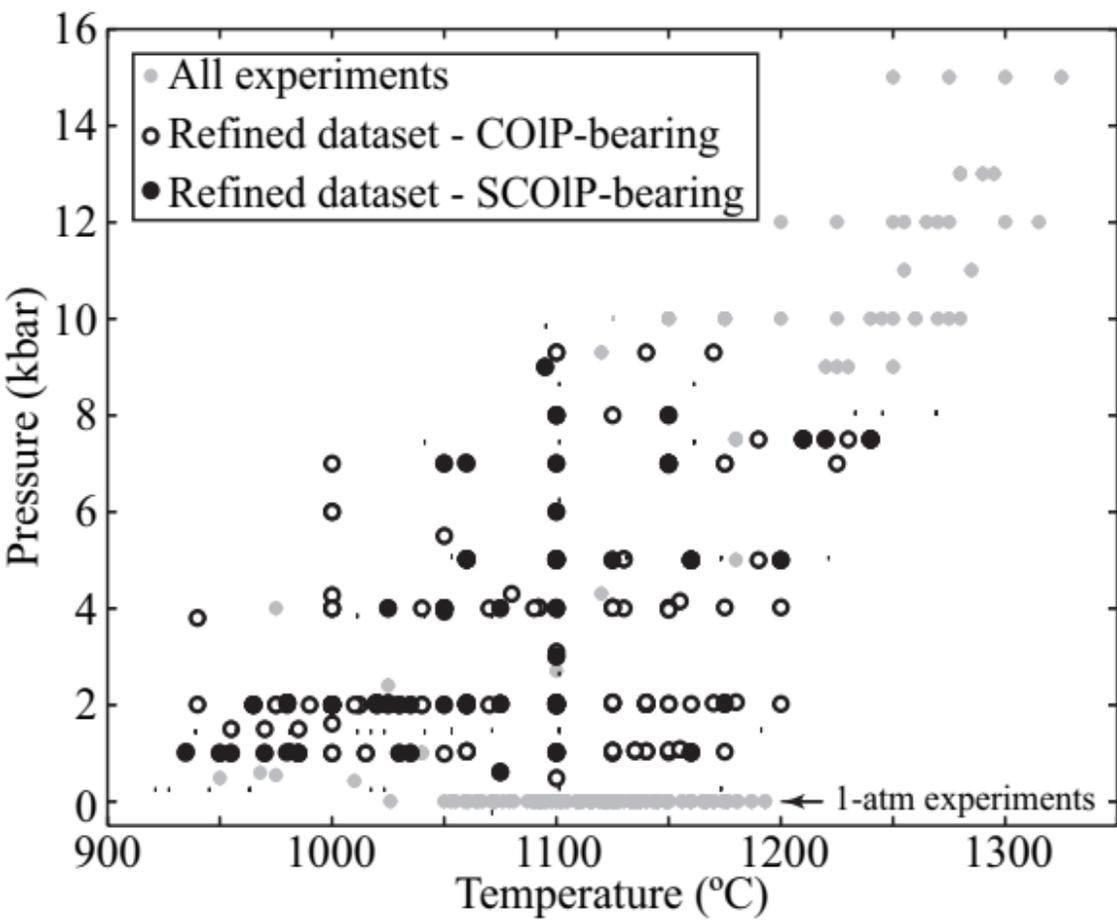


Figure 2

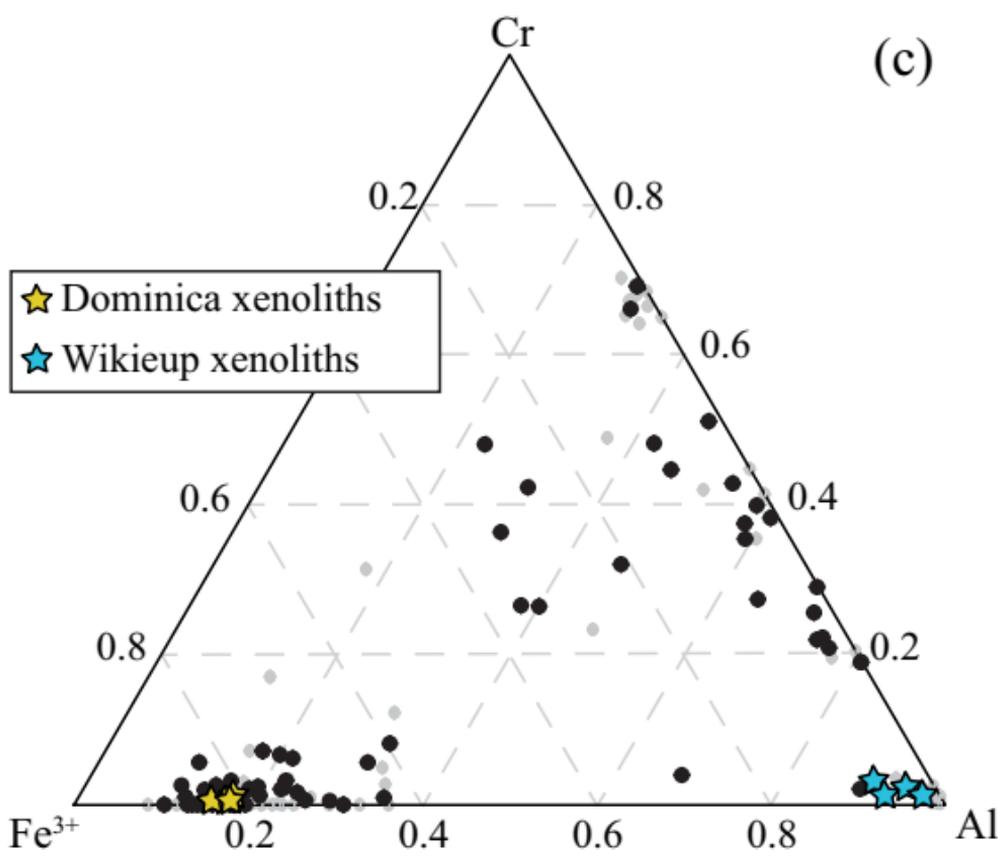
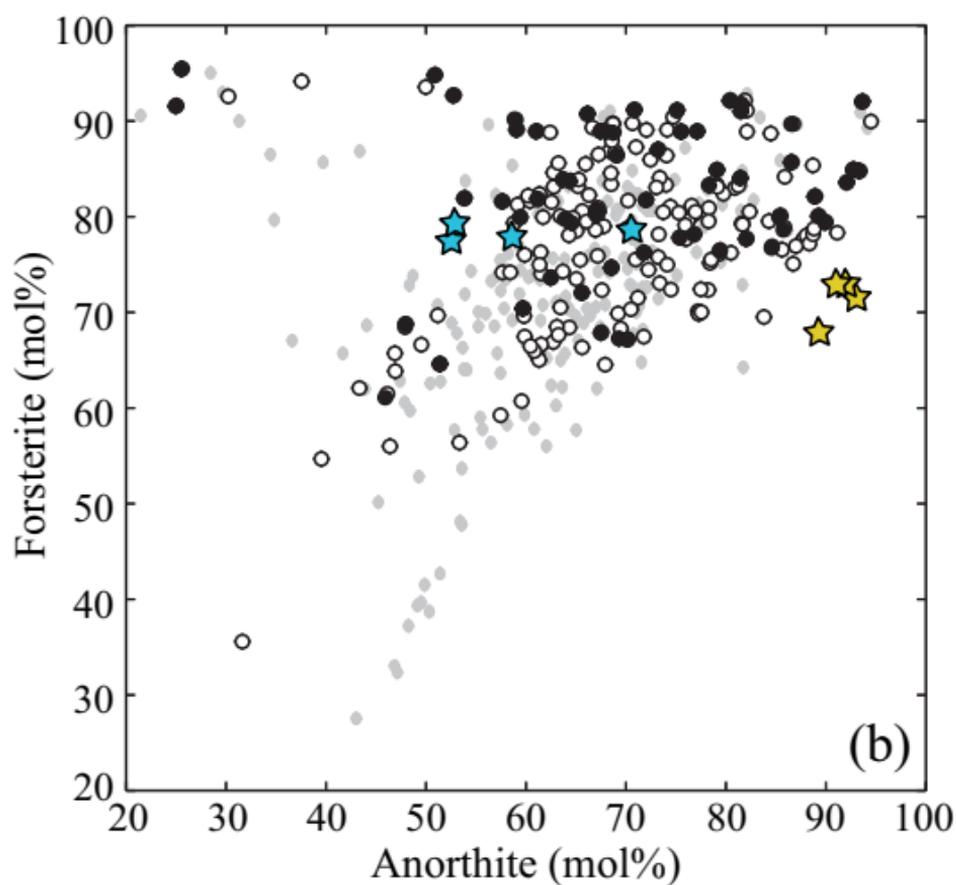
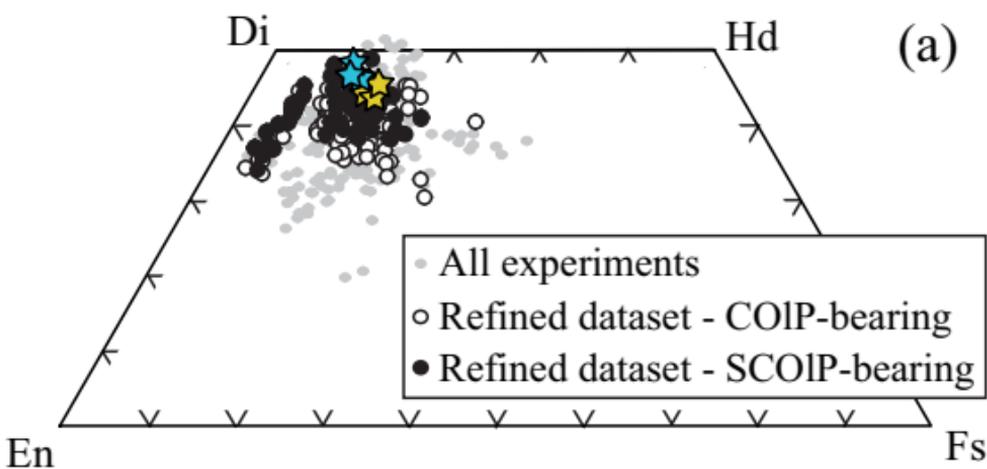


Figure 3

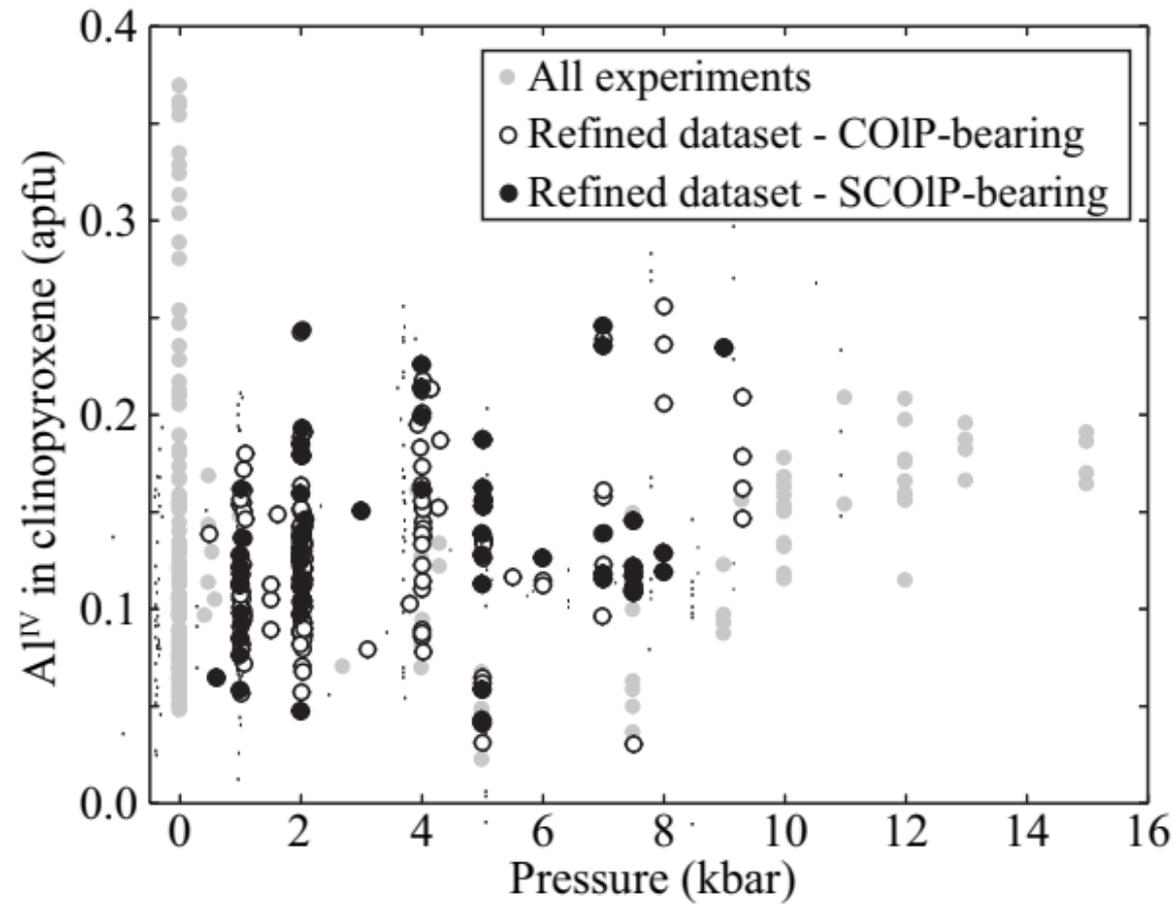
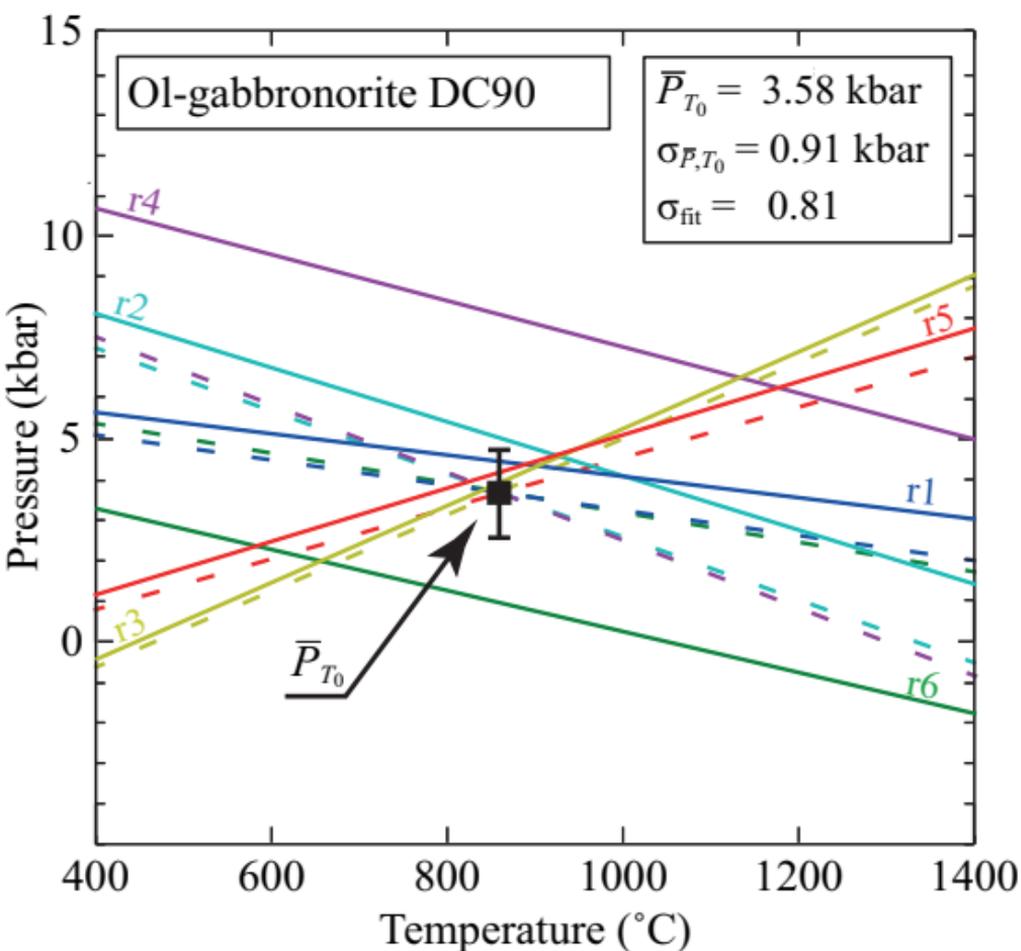


Figure 4



Reactions	P_r (kbar)
r1) fo + an = cenh + cats	4.3 ± 1.4
r2) fa + an = cfs + cats	4.8 ± 1.6
r3) fo + abh = cenh + jd	4.0 ± 1.3
r4) di + an + sp = cenh + 2cats	8.1 ± 4.4
r5) fa + abh + 2sp = cenh + jd + 2herc	4.3 ± 7.2
r6) 2fa + 2abh + mt = 2cfs + 2acm + herc	0.8 ± 2.4

Figure 5

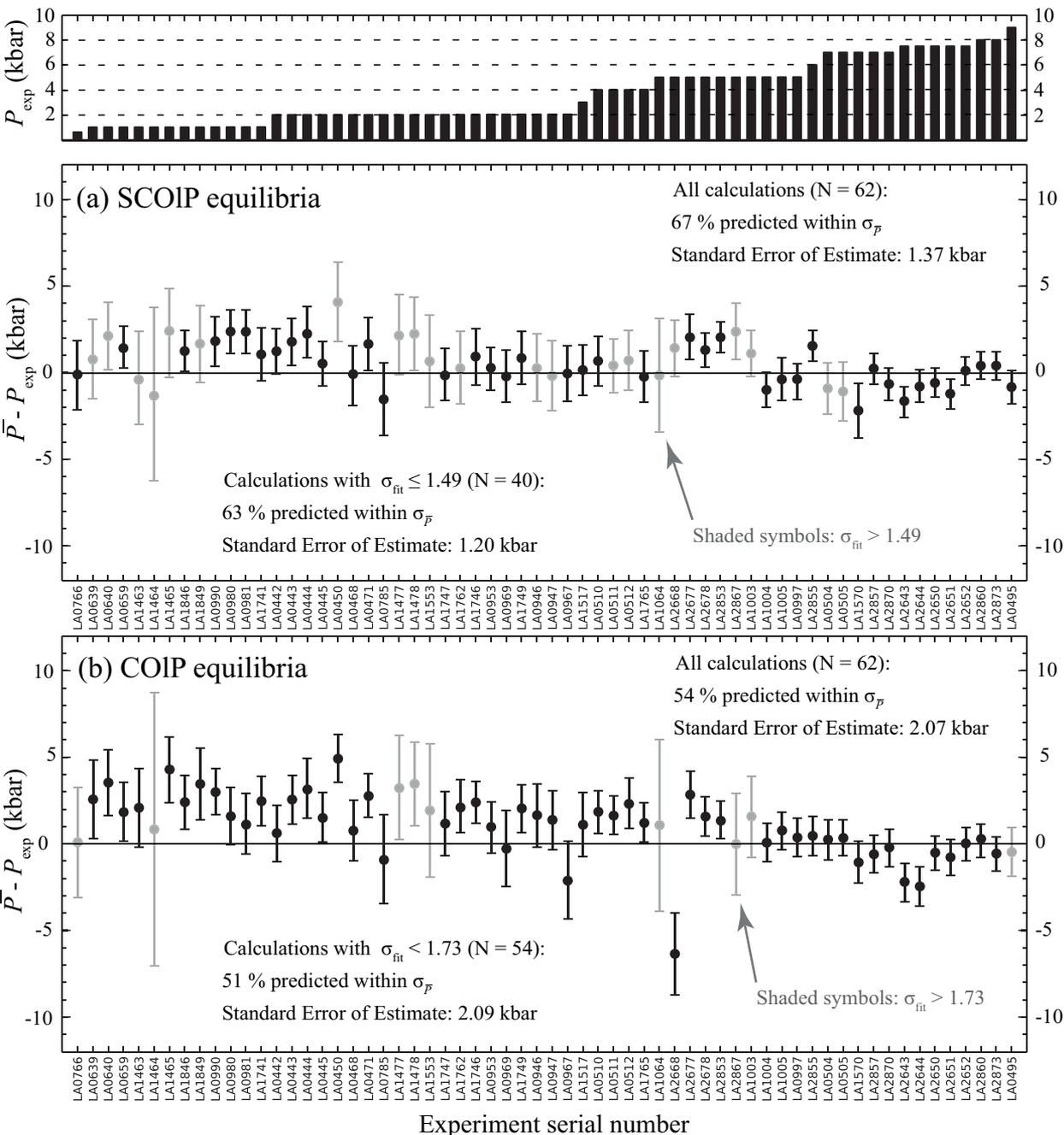


Figure 6

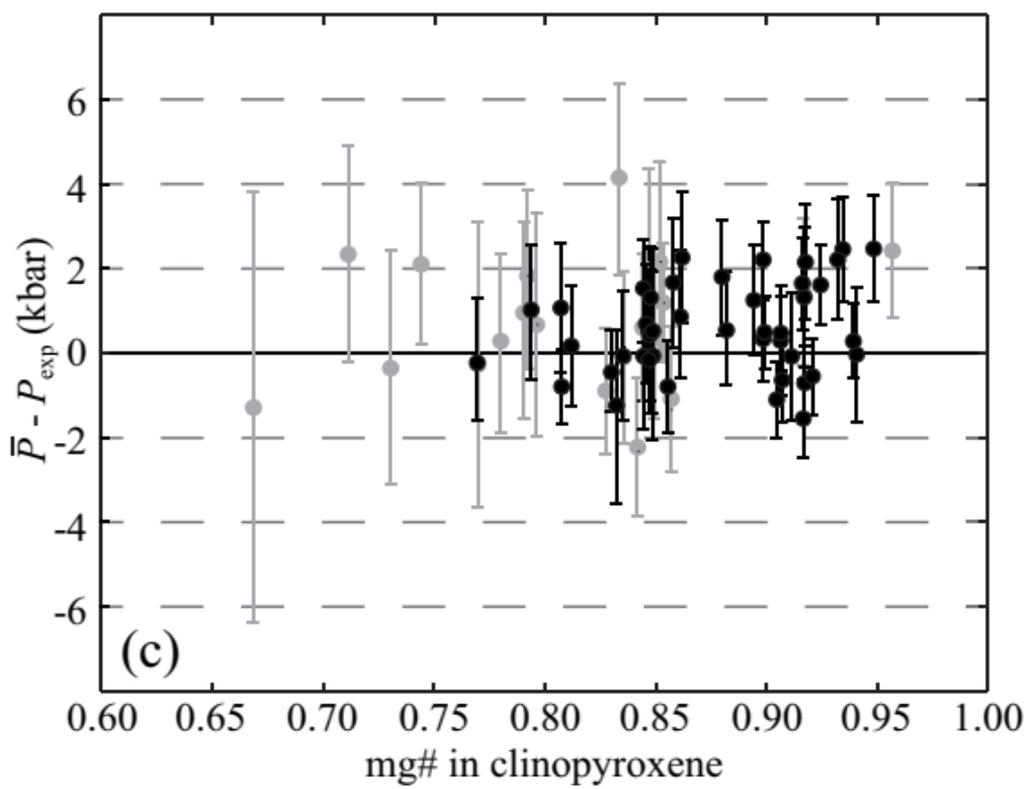
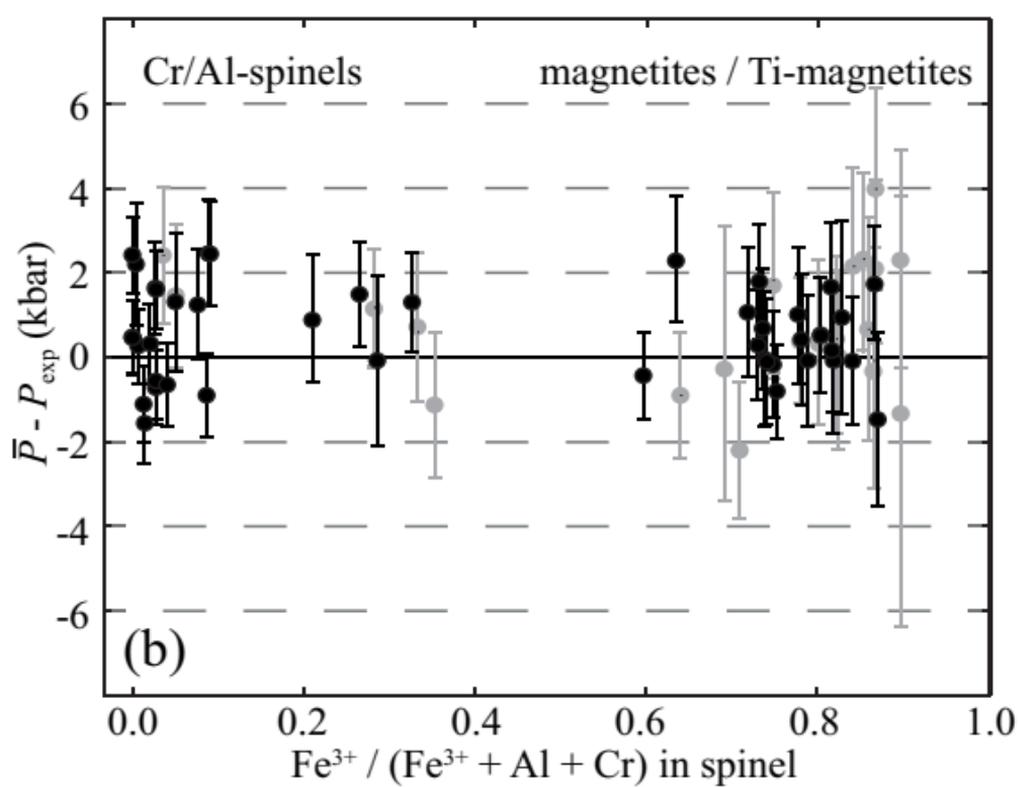
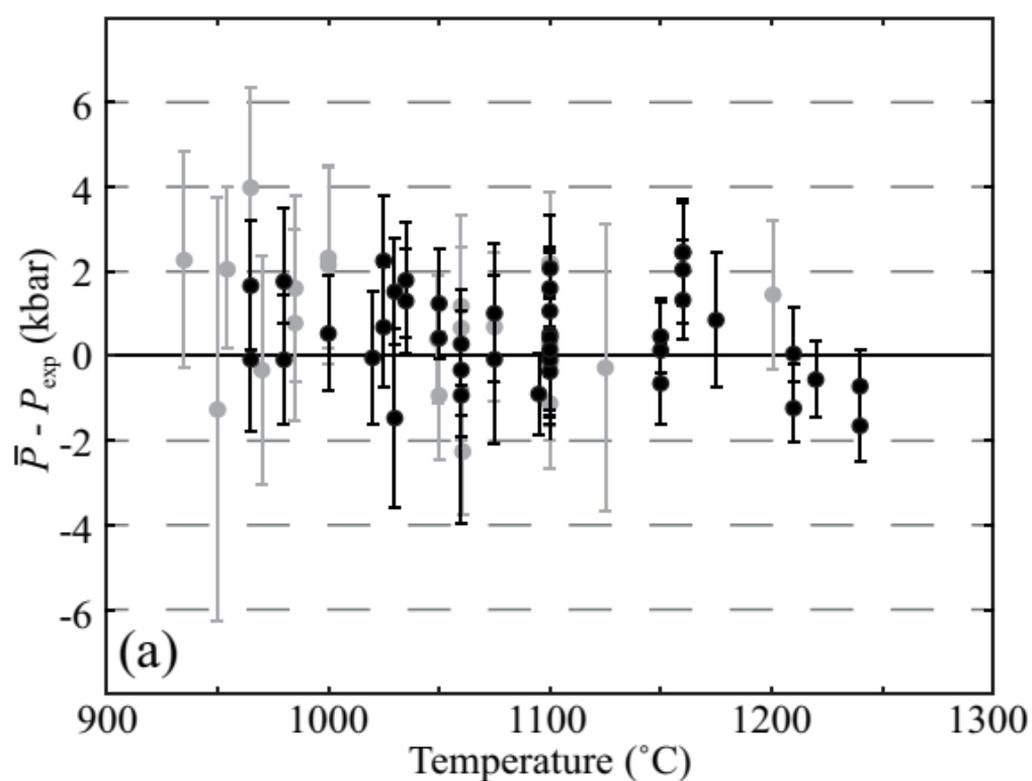


Figure 7

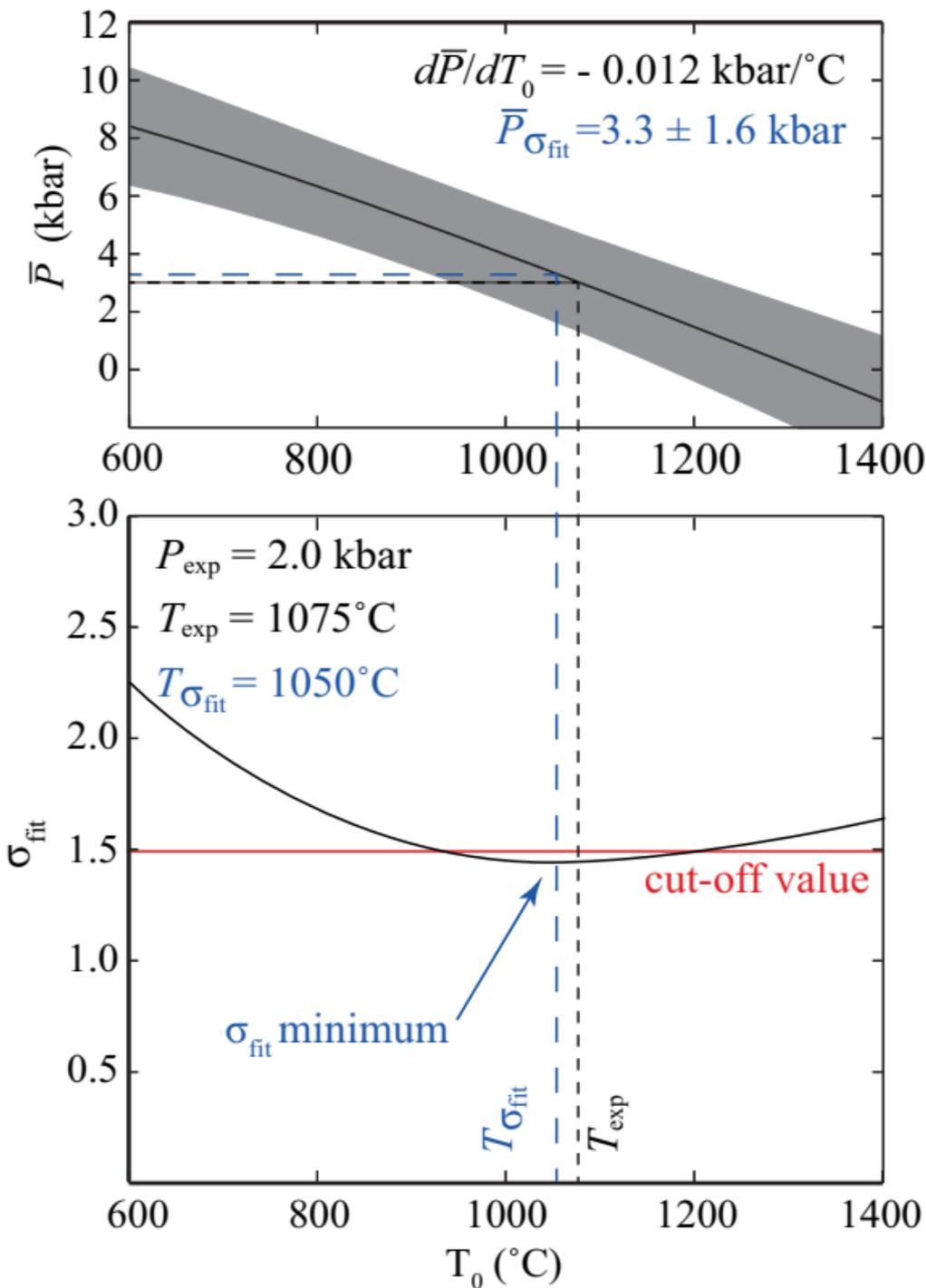


Figure 8

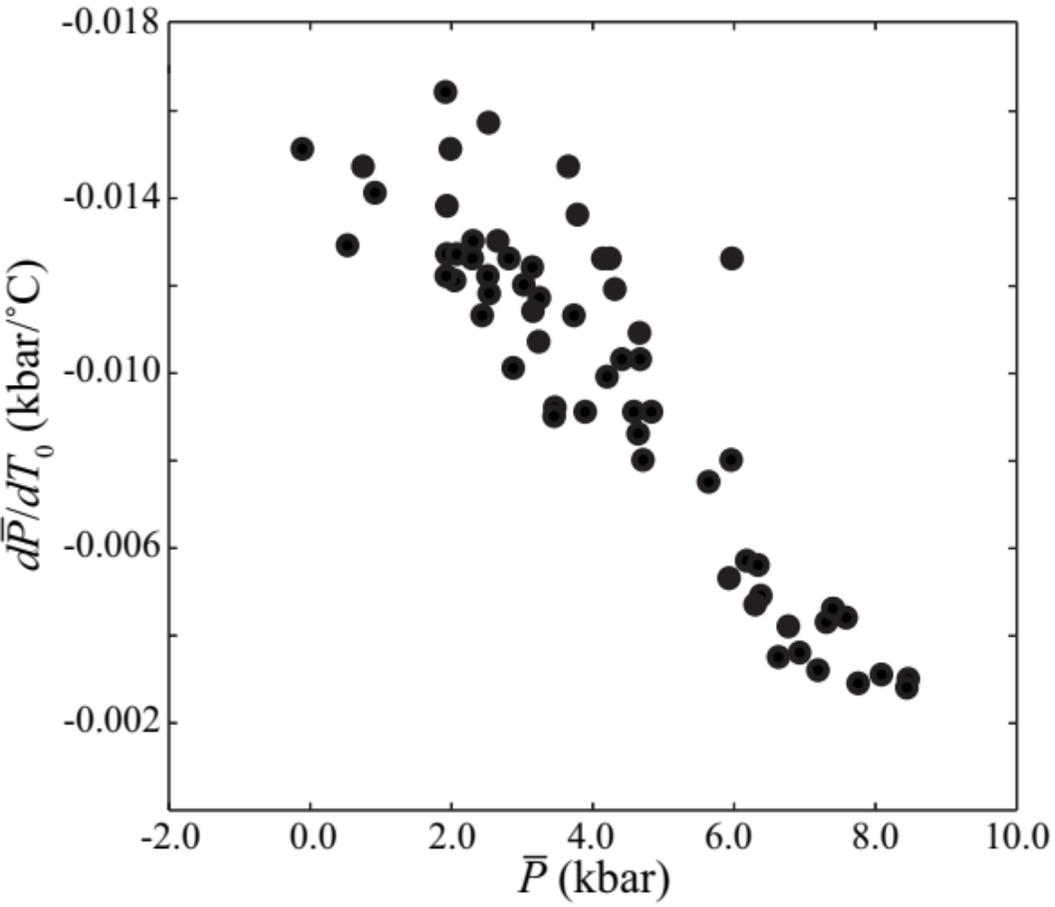


Figure 9

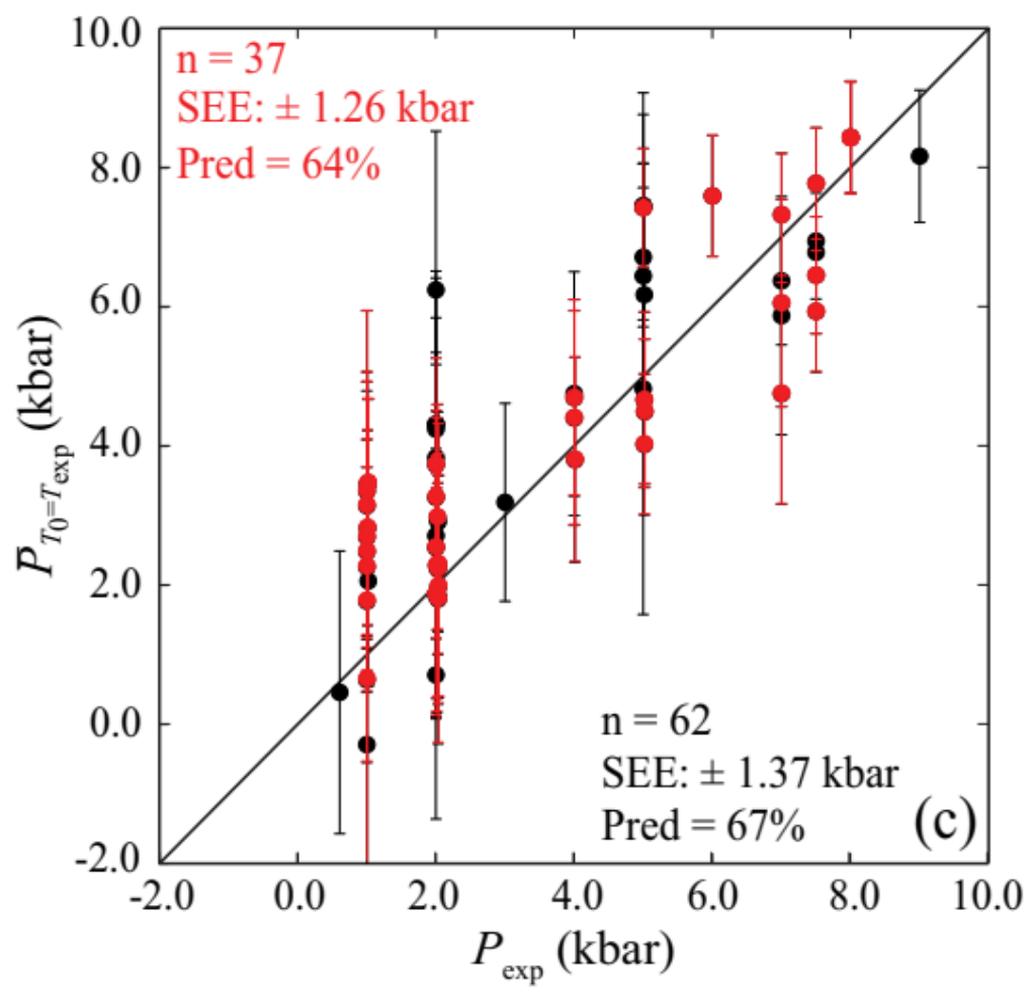
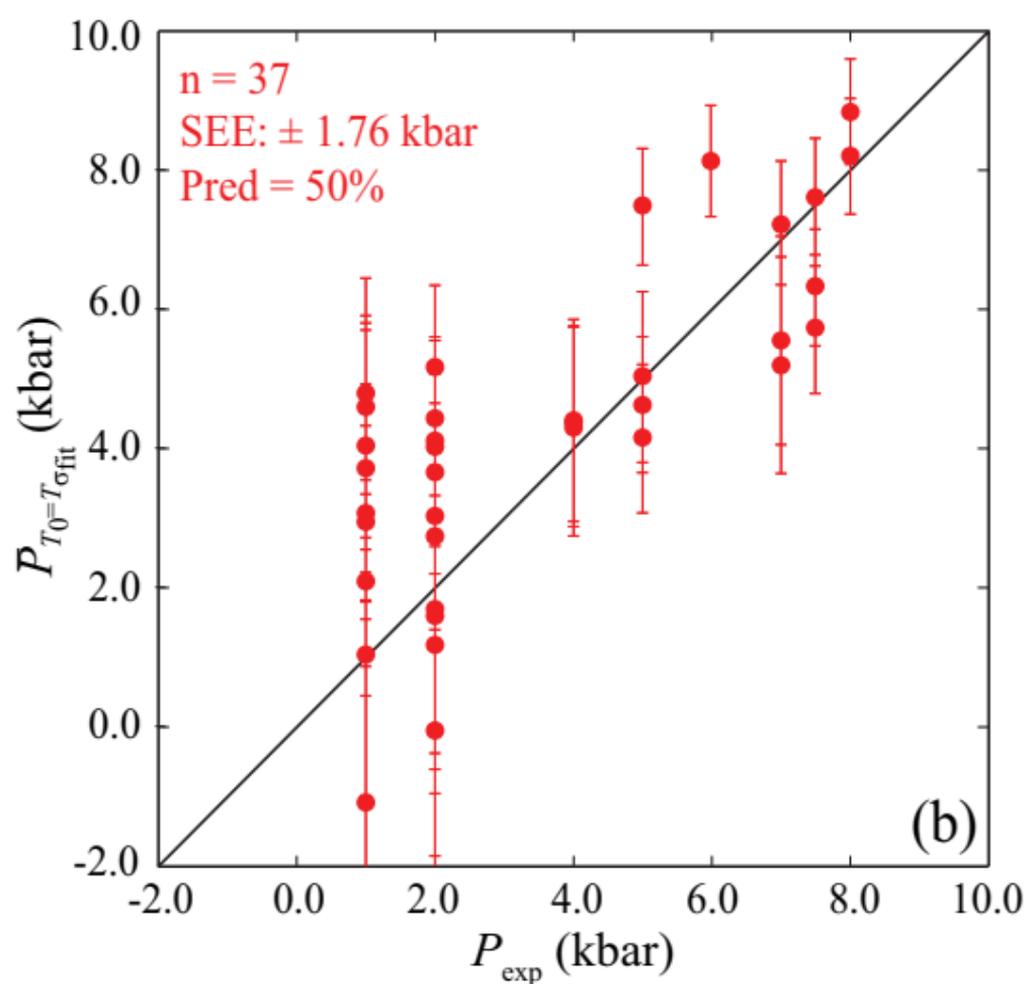
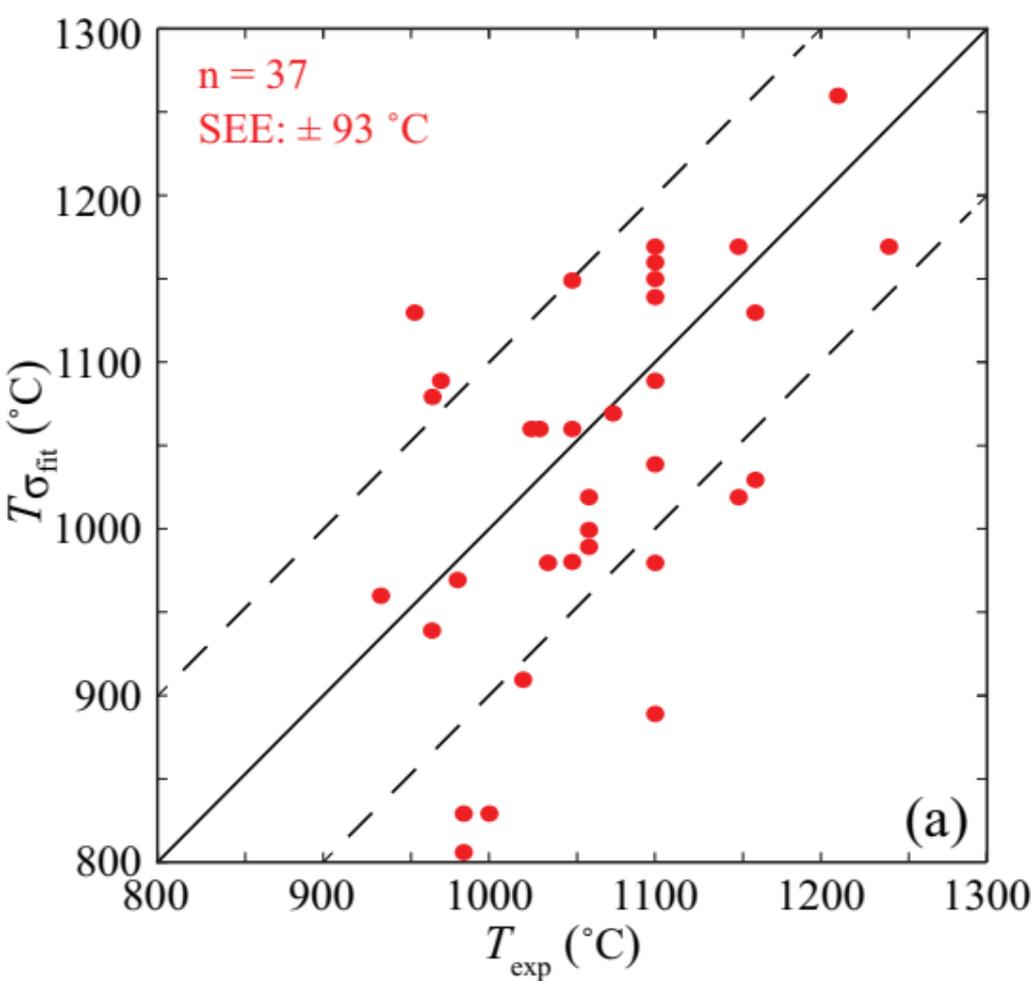


Figure 10

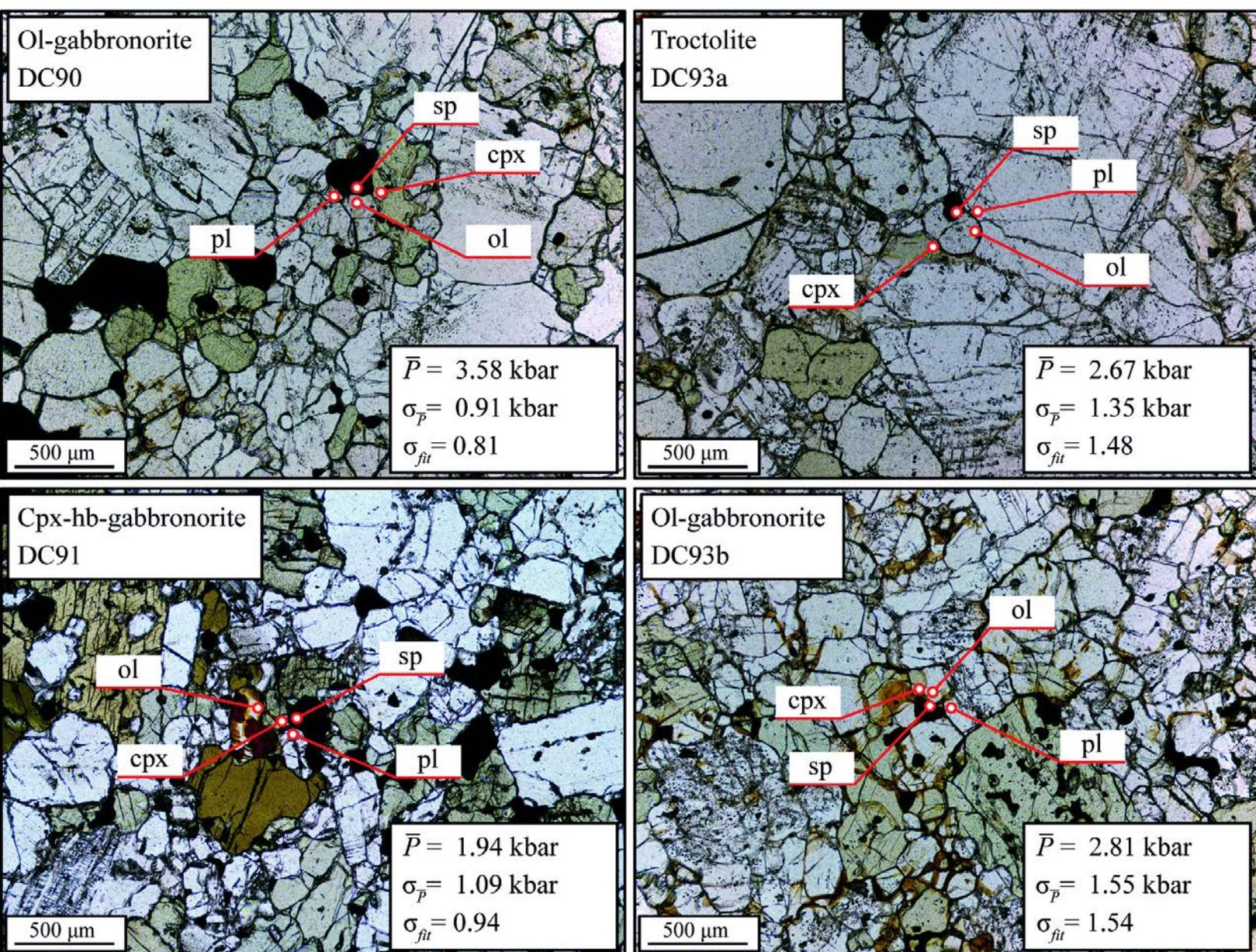


Figure 11

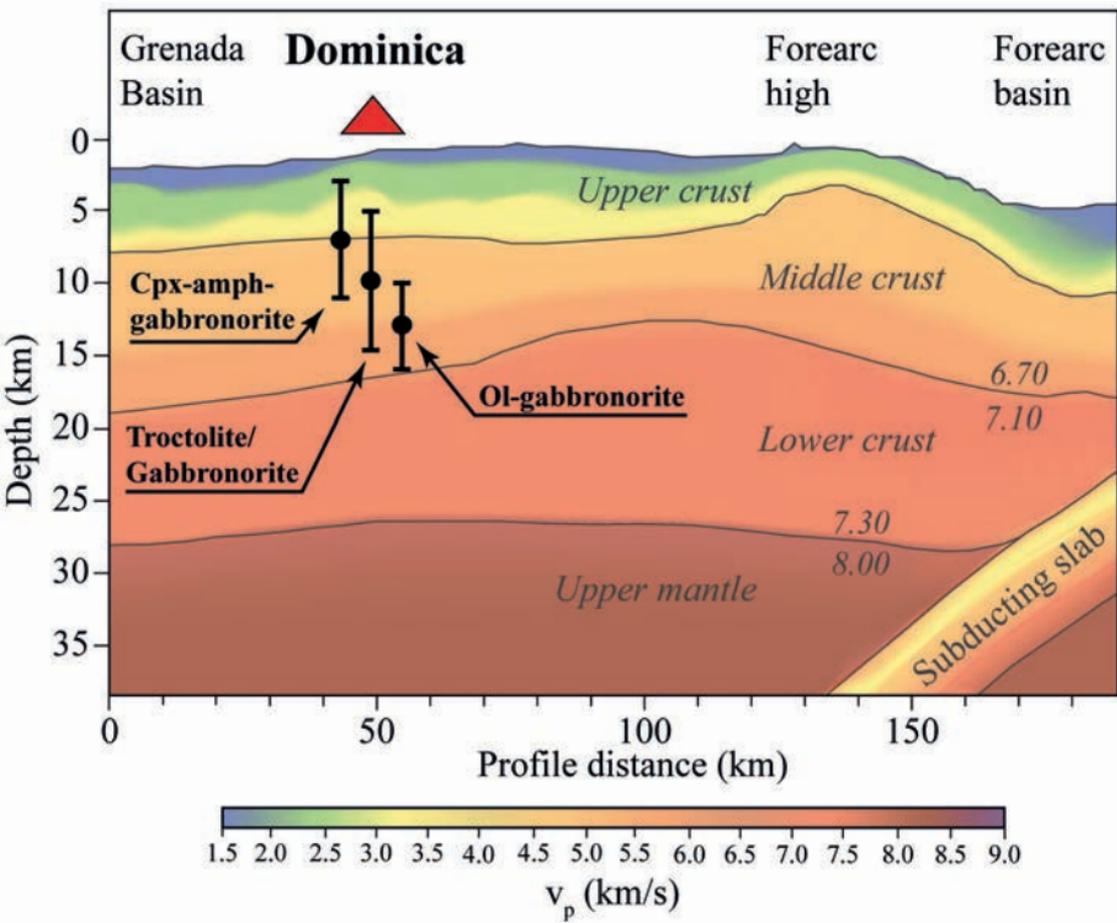


Figure 12

